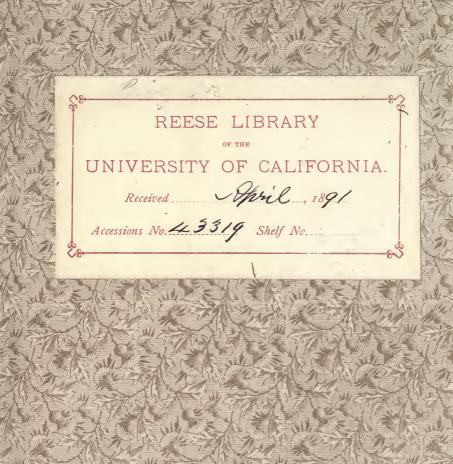


THE CONTINUOUS GIRDER

HOWE

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THE THEORY

OF

THE CONTINUOUS GIRDER:

Its Application to Girders with and without Variable Cross-sections.

BY

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PREFACE.

Continuous bridges—with the exception of draw bridges—are not considered economical and are not designed by American engineers. This probably accounts for the brief treatment the theory of the continuous girder receives in text books and engineering literature.

With one or two exceptions, all American treatises consider the moment of inertia as constant and deduce two equations for the moment over any support, one to be applied when the loads are on the left of the support and the other when the loads are on the right. These two equations combined would, of course, give the moment over any support for any load, but only when the moment of inertia could be assumed as constant, as in girders with parallel flanges, where it might be undue refinement to consider the cross-section as variable.

American engineers have, until recently, assumed the moment of inertia as constant even when the flanges of the girder were considerably inclined, as in the Sabula draw bridge, relying upon the factor of ignorance to cover all discrepancies which might arise from the assumption. A few years ago the modulous of elasticity was quite variable in large bridges—though all engineers considered it as constant in computations—but now, since experience shows that iron and steel can be manufactured with, practically, a constant modulus of elasticity, it may safely be considered as constant, without leading to any appreciable error.

Although no material may be saved by considering the

moment of inertia as variable yet, if it is so considered, the material will be placed where it will do the best service.

Girders with inclined chords, computed and designed as if their cross-sections were constant, have more material than is necessary in some of the compression pieces, and not enough in some of the tension members, as is clearly shown by the results on page 77 of the text.

The object of the following pages is to present to computers, engineers and students a complete mathematical treatment of the theory of the continuous girder, and show how it can be applied to any girder—especially to fixed girders and girders of two and three spans—under any conditions.

With the single assumption that the modulous of elasticity is constant a general equation has been deduced for the moment over any support of any girder under any conditions of loading, of any length of spans, of variable or constant cross-section and for supports at any level.—By difference of level of the supports is meant any change of level which may take place when the girder is in position. It is evident that such a change alone would affect the moments.—From this equation special equations are readily deduced for any particular case. To illustrate the simplicity of the transformations necessary for any special case and also for the convenience of engineers, equations have been given for all the special cases usually discussed in text books, and also equations for these cases when the moment of inertia is considered as variable.

The usual general equations for reactions, deflections and intermediate moments and their transformed equations for special cases are also presented.

Several examples are solved to illustrate the application of the formulas and to show that the processes are almost mechanical when the formulas are thoroughly understood. From those examples which are solved considering the moment of inertia as constant and then variable, a good idea of the manner in which the moment of inertia affects the

results can be obtained. In the case of the Sabula draw there is a difference of about twenty per cent.

In all pin connected bridges the loads are considered to be concentrated at the apices or panel points. In computing the moments for such girders Table I. will be found to be very convenient, as in it are found expressions for

 $2k-3k^2+k^3$ and $k-k^3$ for all values of $k=\frac{a}{l}$ from 0.001 to 0.999 inclusive.

The works enumerated under References, page 107, have been consulted, and some of their parts used without any material change, for which credit is given in foot notes.

The author is indebted to R. H. Brown, C. E., and Geo. H. Hutchinson, C. E., for valuable assistance and suggestions.

M. A. H.

TERRE HAUTE, IND., September, 1888.



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NOMENCLATURE.

r=The number of the support just at the *left* of the $r^{\prime h}$ span.

 l_r =The horizontal length of the span r.

 P_r =Any concentrated load in the r^n span.

 w_r =Any uniform load, per lineal foot, in the $r^{\prime h}$ span.

 a_r =The distance from the left support r to any concentrated load P_r . a_r = $k_r l_r$.

 a_r' =The distance from the support r to the point where the uniform load ends in the r' span.

 a_r'' =The distance from the left support r to the point where the uniform load begins in the r'^h span.

$$k_r = \frac{a_r}{l_r}$$
, or, $a_r = k_r l_r$.

 x_r =The distance from the left support r to any point in the r^{th} span.

 M_r =The bending moment over the support r.

 M_m =The bending moment over any support m.

 M_x =The bending moment at any section, x_r from the support r.

M.=The bending moment at the center of any span of a girder.

 S_r =The shear just at the *right* of the support r.

 S_r' =The shear just at the *left* of the support r.

 R_r =The reaction at the support r, and equals S_r+S_r .

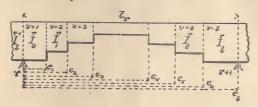
 h_r =The distance the support r is below some specified horizontal line of reference.

 Σ =Sign of summation.

 t_r =Tangent of the angle that the elastic line makes with the horizontal over the support r.

 y_r =The deflection of the girder at any section x_r ; y_r is measured from the horizonal line of reference.

s=The number of spans.



e_i=The distance from the left support r to the point where the moment of inertia of the section of the girder changes for the first time in the span r.

e₂=The distance to the point where it changes the second

time.

 e_v =The distance to any point where the moment of inertia changes, always measured from the left support r and in the span r.

 I_t =The moment of inertia between the last value of e_v and the end of the r^{th} span. See Fig.

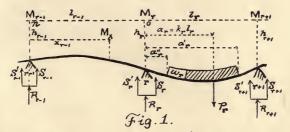
The moment of inertia between the support r and the point e_r .

 \vec{I}_i =The moment of inertia between e_i and e_i in the r^{th} span.

 I_x =The moment of inertia at any section in the r^n span.

E=The modulus of elasticity.

GENERAL RELATIONS.



* In the r^h span of a continuous girder, whose length is l_r , Fig. 1, take a point o vertically above the r^h support, as the origin of co-ordinates, and the horizontal through o as the axis of abscissas. Suppose any section of the girder at a distance x_r from the left support r, and between this section and the support r a weight r distant from r, $a_r = k_r l_r$.

Now, if the girder is continuous over any number of supports there will be a bending moment over each of them (if the ends are not fixed the bending moment over the first and last support will equal zero): M_{r-1} over r-1, M_r over r, M_{r+1} over r+1, etc., also, there will be a shear S_r just at the right and S_{r+1} just at the left of r, S_{r+1} , just at the right and S_{r+1} just at the left of r+1, etc. If there is equilibrum the following conditions must be fulfilled:

^{*} See "Continuirlichen Und Einfachen Träger," page 4, by Prof. Weyrauch. "Theory and Calculation of Continuous Bridges," page 52, by Prof. Merriman.

[&]quot;Strains in Framed Structures," page 244, by Prof. DuBois.

at that support, or

The algebraic sum of all the horizontal forces must be zero. I. The algebraic sum of all the vertical forces must be zero. II.The algebraic sum of the moments of all the forces must III.be zero. From the third condition we have for any section in the rth span, $M = M_r + S_r x_r - P_r (x_r - a_r) \dots x_r = a_r \dots (1)$ If in (1) we make $x_r = l_r$, $M_r = M_{r+1}$, and it becomes $M_{r+1} = M_r + S_r l_r - P_r (l_r - a_r)$, or, since $a_r = k_r l_r$, we have From (2). $S_r = \frac{M_{r+1} - M_r}{L} + P_r (1 - k_r) \qquad (3)$ And if there is no load in the span r this becomes $S_r = \frac{M_{r+1} - M_r}{I} \qquad (4)$ $S'_{r+1} = P_r - S_r$, therefore from (3), $S'_{r+i} = \frac{M_r - M_{r+i}}{l} + P_r k_r$, or better, And if there is no load in the span r-1, (5) becomes The reaction at any support equals the sum of the shears

The above formulas were deduced under the supposition that there was but a single concentrated load P_r in the span r. If there be more than one concentration, the formulas become:

(1)
$$M_r = M_r + S_r x_r - \Sigma P_r (x_r - a_r) \dots x_r = a_r \dots (8)$$

(2)
$$M_{r+j} = M_r + S_r l_r - \Sigma P_r l_r (1-k_r)$$
 (9)

Any partial uniform load in span r.-If w, represents the uniform load per lineal foot in the span r, we can write

$$\Sigma P_r = \int_{a_r=a_r'}^{a_r=a_r'} w_r da_r$$
 or, since $a_r = k_r l_r$

$$\sum_{a_r=a_r''}^{a_r=a_r''} w_r da_r$$
 or, since $a_r = k_r l_r$

$$\sum_{a_r=a_r''}^{a_r=a_r''} v_r da_r$$

$$\sum_{a_r=a_r''}^{a_r=a_r''} v_r da_r$$



The last expression in (12) indicates the difference of the values of the parenthesis when k_r equals $\frac{a_r}{I}$ and $\frac{a_r}{I}$ respectively.

Substituting (12) in the above equations, we have,

From (8),

$$M_{r}^{r} = M_{r} + S_{r} x_{r} - \int_{a_{r}^{r}}^{a_{r}^{r}} w_{r} l_{r} dk_{r} (x_{r} - k_{r} l_{r}) \dots x_{r} = a_{r}^{r} \dots (13)$$

From (9),

Or

$$M_{r+1} = M_r + S_r \ l_r - w_r \ l_r^2 \ (k_r - \frac{1}{2} \ k_r^2) \frac{k_r - \frac{a_r'}{l_r}}{k_r - \frac{a_r''}{l_r}} (15)$$

From (10),

From (11),

$$S_{r}^{\prime} = \frac{M_{r-i} - M_{r}}{l_{r-i}} + \frac{1}{2} w_{r-i} l_{r-i} k_{r-i}^{2} = \frac{a_{r-i}^{\prime}}{l_{r-i}} \dots \dots (17)$$

$$k_{r-i} = \frac{a_{r-i}^{\prime}}{l_{r-i}} \dots \dots \dots \dots (17)$$

Uniform load over entire span r.—If the entire span is uniformly loaded, we have $a_r' = l_r$, $a_r'' = o$ and $a_r = \frac{1}{2} l_r$.

From (13), (In this case $a'_r = x_r$, $a''_r = o$ and $a_r = \frac{1}{2} x_r$ since the load considered cannot extend beyond x_r).

From (15),

$$M_{r+i} = M_r + S_r \ l_r - \frac{1}{9} \ w_r \ l_r^2 \ \dots \ \dots \ (19)$$

From (16),

$$S_r = \frac{M_{r+1} - M_r}{l_r} + \frac{1}{2} w_r l_r \dots$$
 (20)

From (17),

These twenty-one general equations apply to all spans and conditions of loading, but before they can be solved, it will be necessary to find expressions for M or S in terms of the loads, lengths of the spans, etc.

This can be done by introducing the equation of the elastic line and the theorem of three moments. The following expression is obtained for the bending moment over any support m; E alone, being constant.

$$*M_{m} = \frac{-c_{m}\beta_{2}\beta_{3}...\beta_{m-l}}{(d_{s+l}\beta_{l}'')\beta_{2}''...\beta_{m-l}''} \sum_{r=s}^{r=m} \left\{ (A_{r}+Y_{r}+X_{r}'+X_{r-l}'')d_{s-r+2}+B_{r}d_{s-r+l}' \right\} + \frac{-d_{s-m+2}\beta_{s-l}''\beta_{s-l}'\beta_{s-2}''...\beta_{m}''}{(c_{s+l}\beta_{s})\beta_{s-l}'\beta_{s-l}'\beta_{s-2}...\beta_{m}'} \sum_{r=l}^{r=l} \left\{ (A_{r}+Y_{r}+X_{r}'+X_{r-l}'')c_{r}+B_{r}c_{r+l} \right\} (A)$$

$$|\beta_{s-1}'' = |\beta_m''|, \quad |\beta_{s-1} = |\beta_m, \quad |\beta_{s-1} = |\beta_{m-1}, \quad |\beta_{s-1}'' = |\beta_{m-1}''|.$$

In which the respective expressions have the following values:

^{*} A complete demonstration of this equation is given in the Appendix and should be thoroughly understood before any attempt is made to apply equation A practically.

Ar

Concentrated loads-

$$A_r = -\Sigma P_r l_r^2 (2 k_r - 3 k_r^2 + k_r^3) \theta_{r-1} \dots \text{See} (57) \dots (i)$$

Partial uniform load -

From (12),

$$k_r = \frac{a_r'}{l_r}$$
 therefore (i) becomes $k_r = \frac{a_r'}{l_r}$

$$A_{r} = -w_{r} l_{r}^{3} \left\{ k_{r}^{3} - k_{r}^{3} + \frac{k_{r}^{4}}{4} \right\} \begin{pmatrix} a'_{r} = k_{r} l_{r} \\ \theta_{r-1} \\ a''_{r} = k_{r} l_{r} \end{pmatrix}$$
(74)

Uniform load over all-

$$A_r = -\frac{1}{4} w_r l_r^u \theta_{r-1} \dots (75)$$

 \boldsymbol{B} .

Concentrated loads-

$$B_r = -\Sigma P_r l_r^2 (k_r - k_r^3) \theta_{r+1} \dots$$
 See (58) (j)

Partial uniform loads-

$$B_{r} = -w_{r} l_{r}^{s} \left\{ \frac{2 k_{r}^{s} - k_{r}^{s}}{4} \right\} \begin{cases} a_{r}^{c} = k_{r} l_{r} \\ \theta_{r+s} \\ a_{r}^{c} = k_{r} l_{r} \end{cases}$$
 (76)

Uniform load over all-

$$B_r = -\frac{1}{4} w_r \, \tilde{\ell}_r^s \, \theta_{r+i} \, \dots \, (77)$$

X'_{r}

Concentrated loads-

$$X_r = -F_r \stackrel{\Sigma}{=} P_r l_r (1-k_r) \theta_{r-1} + H_r \theta_{r-1} \dots \dots \dots (h)$$

Partial uniform load—

$$X'_{r} = -F'_{r} w_{r} l_{r}^{2} \theta_{r-i} \left\{ k_{r} - \frac{1}{2} k_{r}^{2} \right\} a'_{r} = k_{r} l_{r} + H_{r} \theta_{r-i} . . . (78)$$

Uniform load over all-

$$X'_r = -\frac{1}{2} F'_r w_r l_r^2 \theta_{r-t} + H_r \theta_{r-t} \dots \dots \dots (79)$$

X''_{r-1}

Concentrated loads-

$$X_{r-i}^{"} = -2 F_{r-i}^{"} \Sigma P_{r-i} l_{r-i} (1-k_{r-i}) \theta_r - H_{r-i} \theta_r \dots \text{See} (59) \dots (h)$$

Partial uniform load-

$$X_{r-i}^{"} = -2 F_{r-i}^{"} \theta_r w_{r-i} l_{r-i}^2 \left\{ k_{r-i} - \frac{1}{2} k_{r-i}^2 \right\} \begin{pmatrix} a'_{r-i} = k_{r-i} l_{r-i} \\ -H'_{r-i} \theta_r \end{pmatrix}. (80)$$

Uniform load over all-

H_{r}

Concentrated loads-

$$r=1$$

$$H_r = \sum_{v} \bigwedge_{l_r} \left\{ \frac{\sum P_r \left(e_v - a_r \right)^3}{l_r} + \frac{3 \left(l_r - e_v \right)}{l_r} \sum P_r \left(e_v - a_r \right)^2 \right\} \dots (f)$$

Partial uniform load-

Substitute the following expressions in (f):

$$\sum_{r} P_{r} (e_{r} - a_{r})^{3} = \int_{a_{r}}^{c} w_{r} d a_{r} (e_{r} - a_{r})^{3} \dots a_{r} \leq e_{r} \dots (82)$$

$$\sum_{\alpha_r} P_r (e_v - a_r)^2 = \int_{\alpha_r'}^{\alpha_r'} w_r \, d \, a_r \, (e_v - a_r)^2 \, \dots \, d_r' \leq e_v \, \dots \, . \quad (83)$$

Probably the easiest way to handle this case is to consider the uniform load as extending from the left support to the farther end of the load and integrate between the limits $a'_r \leq e_v$ and $a''_r = o$; then consider the load as extending from the left support to the nearer end of the load and integrate between the limits $a'_r = a''_r \leq e_v$ and o, and take the difference of the results.

Uniform load over all-

H'_{i}

Concentrated loads-

$$H'_{r} = \sum_{v=l_{r}} \sum_{r}^{r} \left\{ \frac{\sum P_{r} (c_{v} - a_{r})^{3}}{l_{r}} - \frac{3 e_{v}}{l_{r}} \sum P_{r} (e_{v} - a_{r})^{2} \right\} \dots (g)$$

Partial uniform load-

$$\Sigma P_r (e_v - a_r)^3 = \int_{a_r}^{a_r'} w_r d a_r (e_v - a_r)^3 \dots a_r' \leq e_v \dots (82)$$

$$\Sigma P_r (e_v - a_r)^2 = \int_{a_r'}^{a_r'} w_r d a_r (e_v - a_r)^2 \dots a_r' \leq e_v \dots$$
 (83)

Uniform load over all-

$$\sum_{r} P_{r} (e_{r} - a_{r})^{3} = \int_{0}^{e_{r}} w_{r} d a_{r} (e_{r} - a_{r})^{3} = \frac{1}{4} w_{r} \epsilon_{r}^{4} \dots \dots (84)$$

$$\Sigma P_{r} (e_{v} - a_{r})^{s} = \int_{0}^{e_{v}} w_{r} d a_{r} (e_{v} - a_{r})^{s} = \frac{1}{\beta} w_{r} e_{v}^{s} \dots (85)$$

For all loads-

$$v=I F_r = \sum_{r} \bigwedge_{r} \left\{ c_r \left(3 - \frac{3 e_r}{l_r} + \frac{c_r^2}{l_r^2} \right) = \frac{l_r^3 - (l_r - e_r)^3}{l_r^2} \right\} \text{See (48)} . . . (b) v=l_r$$

$$v=1$$

$$F_r = \sum_{v=l_r} \stackrel{?}{\underset{v}{\triangle}} e_v^2 \left\{ \frac{3}{l_r} - \frac{2e_v}{l_r^2} \right\} \dots \text{See (49)} \dots (c)$$

$$v=1$$

$$F_r''=\Sigma \bigwedge_{\stackrel{\sim}{r}} \stackrel{c^3}{\sim} e_v^3 \frac{1}{l_r^2} \dots \dots \dots \text{See (50)} \dots \dots (d)$$

$$v=l_r$$

$$d_{s}=0.$$
 $d_{2}=1.$
$$d_{m}=-2 d_{m-1} \frac{\beta'_{s-m+3}}{\beta'_{s-m+2}} - d_{m-2} \frac{\beta_{s-m+3}}{\beta'_{s-m+2}} . . . See (68) (r)$$

By means of the above equations we can deduce the bending moment over any support of any continuous girder under the single assumption that the modulus of elasticity, E, shall be constant.

E AND I CONSTANT.

If the moment of inertia as well as the modulus of elasticity is constant our equations will be much more simple.

By inspection, we see that (a) equals zero, and hence F_r , F_r' and F_r'' equal zero and also H_r and H_r' ; hence X_r' and X_{r-r}'' equal zero. By reduction we obtain

$$M_{m} = \frac{-c_{m}}{d_{s+i}} \sum_{l_{r}=s}^{r=m} \left\{ (A_{r} + Y_{r}) \ d_{s-r+2} + B_{r} \ d_{s-r+i} \right\}$$

$$+ \frac{-d_{s-m+2}}{c_{s+i}} \sum_{l_{s}}^{r=i} \left\{ (A_{r} + Y_{r}) \ c_{r} + B_{r} \ c_{r+i} \right\} (A_{i})$$

$$c_{j} = 0, \quad c_{2} = 1,$$

$$c_{m} = -2 c_{m-j} \frac{l_{m-2} + l_{m-j}}{l_{m-j}} - c_{m-2} \frac{l_{m-2}}{l_{m-j}} \qquad (p_{i})$$

$$d_{j} = 0, \quad d_{2} = 1,$$

$$d_{m} = -2 d_{m-j} \frac{l_{s-m+3} + l_{s-m+2}}{l_{s-m+2}} - d_{m-s} \frac{l_{s-m+3}}{l_{s-m+2}} \qquad (r_{i})$$

$$Y_{r} = -6 E I \left\{ \frac{h_{r} - h_{r-j}}{l_{r-j}} + \frac{h_{r} - h_{r+j}}{l_{r}} \right\} \qquad (l_{i})$$

$$A_{r}$$

$$Concentrated \ loads - A_{r} = -\sum P_{r} l_{r}^{2} (2 k_{r} - 3 k_{r}^{2} + k_{r}^{5}) \qquad (i_{i})$$

$$Partial \ uniform \ load - A_{r} = -w_{r} l_{r}^{3} - \left\{ k_{r}^{2} - l_{r}^{3} + \frac{k_{r}^{4}}{4} \right\} a_{r}^{a} = k_{r} l_{r}$$

$$Uniform \ load \ over \ all - A_{r} = -\frac{1}{4} w_{r} l_{r}^{2} \qquad (87)$$

$$B_{r}$$

$$Concentrated \ loads - B_{r} = -\sum P_{r} l_{r}^{2} (k_{r} - k_{r}^{3}) \qquad (j_{i})$$

$$Partial \ uniform \ load - B_{r} = -w_{r} l_{r}^{3} \left\{ \frac{2}{2} \frac{k_{r}^{2} - k_{r}^{4}}{4} \right\} a_{r}^{a} = k_{r} l_{r}$$

$$Uniform \ load \ over \ all - B_{r} = -w_{r} l_{r}^{3} \left\{ \frac{2}{2} \frac{k_{r}^{2} - k_{r}^{4}}{4} \right\} a_{r}^{a} = k_{r} l_{r}$$

$$Uniform \ load \ over \ all - B_{r} = -\frac{1}{h} w_{r} l_{r}^{3} \qquad (88)$$

E AND I CONSTANT-Spans equal.

If the spans are equal, the equations of the last case reduce to,

$$M_{m} = \frac{-c_{m}}{d_{s+i}} t \sum_{r=s}^{r=m} \left\{ (A_{r} + Y_{r}) d_{s-r+2} + B_{r} d_{s-r+i} \right\}$$

$$+ \frac{-d_{s-m+2}}{c_{s+i}} \sum_{r=m-i}^{r=i} \left\{ (A_{r} + Y_{r}) c_{r} + B_{r} c_{r+i} \right\} ... (A_{2})$$

$$c_{i} = 0. c_{2} = 1.$$

$$c_{m} = -4 c_{m-i} - c_{m-2} ... (p_{2})$$

$$d_{i} = 0. d_{2} = 1.$$

$$d_{m} = -4 d_{m-i} - d_{m-2} ... (r_{2})$$

Concentrated loads-

Partial uniform load-

$$A_{r} = -w_{r} l^{3} \left\{ k_{r}^{2} - k_{r}^{3} + \frac{k_{r}^{3}}{4} \right\} \frac{a'_{r} = k_{r} l}{a''_{r} = k_{r} l} \dots \dots (90)$$

Uniform load over all— .

B_r

Concentrated loads-

Partial uniform load-

Uniform load over all-

$$B_r = -\frac{1}{4} w_r l^3 \dots (93)$$

In case the supports are at the same level, $Y_r=0$ in all equations.

From (p_2) and (r_2) we obtain the following values for c and c:

$$c_{i} = \pm 0 = d_{i}$$
 $c_{7} = -780 = d_{7}$ $c_{2} = +1 = d_{2}$ $c_{8} = +2911 = d_{8}$ $c_{5} = -4 = d_{3}$ $c_{9} = -10864 = d_{9}$ $c_{4} = +15 = d_{4}$ $c_{10} = +40545 = d_{10}$ $c_{5} = -56 = d_{5}$ $c_{11} = -151316 = d_{11}$ $c_{12} = +209 = d_{6}$ $c_{12} = +564719 = d_{12}$

Explanation of Table I— In the co-efficients $(2 k_r - 3 k_r^2 + k_r^3)$ and $(k_r - k_r^3)$, k_r is a fraction and equals $\frac{a_r}{l_r}$, hence, it is immaterial about the actual values of a and l as long as the ratio is known.

The ratio k_r has been assumed to have all possible values from .001 to .999 inclusive, and the respective values of the above co efficients, carefully computed, are given in Table I. A few trials will prove the great utility of this table, and convince the computer that much time and labor can be saved by its use.

SHEAR.

The moments over the supports being determined, the shears can be found from the following formulas, which apply to all cases.

$$S_r = \frac{M_{r+1} - M_r}{l_r} + Q_r \dots \dots \text{See (10)} \dots (B)$$

$$S'_{r} = \frac{M_{r-1} - M_{r}}{l_{r-1}} + Q'_{r-1} \qquad \text{See (11)} \qquad (C)$$
In which Q and Q have the following values:

$$\begin{array}{l} \textbf{Concentrated loads-} \\ Q_{r} = \Sigma P_{r} (1-k_{r}) \qquad (94) \\ Q'_{r} = \Sigma P_{r} k_{r} \qquad (95) \\ \textbf{Partial uniform load-} \\ Q_{r} = w_{r} l_{r} \left\{ \frac{2}{2} \frac{k_{r} - k_{r}^{2}}{2} \right\} \begin{array}{l} a'_{r} = k_{r} l_{r} \\ a''_{r} = k_{r} l_{r} \end{array} \qquad (96)$$

$$\begin{array}{l} \textbf{Partial uniform load over all-} \\ Q'_{r} = w_{r} l_{r} \left\{ \frac{k_{r}^{2}}{2} \right\} \begin{array}{l} a'_{r} = k_{r} l_{r} \\ a''_{r} = k_{r} l_{r} \end{array} \qquad (97)$$

$$\begin{array}{l} \textbf{Uniform load over all-} \\ Q_{r} = \frac{1}{2} w_{r} l_{r} \qquad (99) \\ \textbf{INTERMEDIATE BENDING MOMENTS.} \\ \text{General equations for all possible cases.} \\ M'_{z} = M_{r} + S_{r} x_{r} - L_{r} \qquad See (8) \qquad (D) \\ \text{In which } \textbf{L} \text{ has the following values:} \\ \textbf{Concentrated loads-} \\ L_{r} = \Sigma P_{r} (x_{r} - a_{r}) \qquad a_{r} = a'_{r} \qquad a'_{r} \leq x_{r} \qquad (100) \\ \textbf{Partial uniform loads-} \\ L_{r} = w_{r} \left\{ a_{r}(x_{r} - \frac{a_{r}}{2}) \right\} \begin{array}{l} a_{r} = a'_{r} \qquad a'_{r} \leq x_{r} \qquad (101) \\ a_{r} = a'_{r} \qquad a'_{r} = a'_{r} \qquad a'_{r} \leq x_{r} \qquad (102) \\ \end{array}$$

DEFLECTION.

General formula. E alone constant.

$$y_{r} = h_{r} + t_{r} x_{r} + \frac{1}{6 E I_{x}^{r}} \left\{ 3 M_{r} x_{r}^{s} + S_{r} x_{r}^{3} - \Sigma P_{r}(x_{r} - a_{r})^{s} \right\}$$

$$+ \frac{1}{6 E I_{x}^{r}} \sum_{v=x_{r}}^{v=1} \left(\frac{I_{x}}{I_{v-1}} - \frac{I_{x}}{I_{v}} \right) \left\{ 3 M_{r} e_{v} (2 x_{r} - e_{v}) + S_{r} e_{v}^{2} (3 x_{r} - 2 e_{v}) \right.$$

$$- \Sigma P_{r} (e_{v} - a_{r})^{s} - 3 (x_{r} - e_{v}) \Sigma P_{r} (e_{v} - a_{r})^{2} \right\} . \text{See (39)} . . (E)$$

$$t_{r+i} = \frac{h_{r+i} - h_{r}}{l_{r}} + \frac{1}{6 E I_{i}^{r}} \left\{ M_{r} l_{r} + 2 M_{r+i} l_{r} + \Sigma P_{r} l_{r}^{2} (k_{r} - k_{r}^{3}) \right.$$

$$+ \frac{1}{6 E I_{i}^{r}} \sum_{v=l_{r}}^{v=1} \left(\frac{I_{i}}{I_{v-i}} - \frac{I_{i}}{I_{v}} \right) \left\{ M_{r} e_{v}^{2} (3 - \frac{2 e_{v}}{l_{r}}) + 2 M_{r+i} \frac{e_{v}^{3}}{l_{r}} + \frac{2 e_{v}^{3}}{l_{r}} \Sigma P_{r} l_{r} (1 - k_{r}) + \Sigma P_{r} (e_{v} - a_{r})^{3} - 3 e_{v} \Sigma P_{r} (e_{v} - a_{r})^{2} \right\} . (45)$$

For uniform loads-

Remember that in the terms containing (x-a) and (e-a), x and e must never be less than a, or rather $a \le x$ or e.

E AND I CONSTANT.

If the moment of inertia is constant the preceding equations reduce to,

$$y_{r} = h_{r} + t_{r} x_{r} + \frac{1}{6 E I} \left\{ \beta M_{r} x_{r}^{2} + S_{r} x_{r}^{3} - \Sigma P_{r} (x_{r} - a_{r})^{3} \right\} . . (E_{s})$$

$$t_{r+s} = \frac{h_{r+s} - h_{r}}{l_{r}} + \frac{1}{6 E I} \left\{ M_{r} l_{r} + \beta M_{r+s} l_{r} + \Sigma P_{r} l_{r}^{2} (k_{r} - k_{r}^{3}) \right\} . . . (45_{a})$$

For uniform loads—

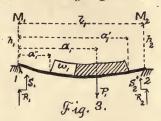
If the supports are at the same level, the terms containing h in the above equations become zero; the equations remain the same in every other particular.

We have now deduced general equations by which we can determine the bending moments, shears and deflections for any continuous girder. In the next chapters we will give numerous examples or special cases illustrating the application of the formulas.

SUPPORTED GIRDERS.

CASE I.

A simple girder resting upon two supports-



(a) E alone constant.

(a) E alone constant.
From (A) we have for M_s and M_s
$M_1 = M_2 = 0 \dots \dots$
From (B) and (C)
$S_t = Q_t \dots \dots$
$S_2 = Q_1 \dots \dots$
Or
$S_i = \Sigma P_i (1-k_i)$. for concentrated loads (106)
$S_{z} = \Sigma P_{s} k_{s}$. for concentrated loads (107)
$S_i=w, \ l_i \ \left\{k_i-\frac{k_i^2}{2}\right\} \begin{array}{c} a_i'=k_i \ l_i \\ \vdots \\ a_i''=k_i \ l_i' \end{array}$ for any uniform load (108)
$S_{2}=w_{i} l_{i} \left\{\frac{k_{i}^{2}}{2}\right\} \stackrel{a'_{i}=k_{i} l_{i}}{\vdots} for any uniform load (109)$
$S_i = \frac{1}{2} w_i l_i \dots$ for uniform load over all (110)

$$S_{z}' = \frac{1}{2} w_{i} l_{i} \dots$$
 for uniform load over all (111)

From (D) the bending moment at any section x_i is

Substituting the values of S_i and L_i we have,

For concentrated loads—

$$M_{x} = \sum P_{i} (1-k_{i}) x_{i} - \sum P_{i} (x_{i}-a_{i}) \dots a \leq x \dots (113)$$

For any uniform load—

$$M_{x}' = x_{i} w_{i} l_{i} \left\{ k_{i} - \frac{k_{i}^{2}}{2} \right\} \begin{array}{c} a_{i}' \\ - w_{i} \\ a_{i}'' \end{array} \left\{ a_{i} \left(x_{i} - \frac{a_{i}}{2} \right) \right\} \begin{array}{c} a_{i} = a_{i}' \\ a_{i} \leq x_{i} \\ a_{i} = a_{i}'' \end{array}$$
(114)

For uniform load over all—

$$M_{x}^{i} = \frac{1}{2} w_{i} l_{i} x_{i} - \frac{1}{2} w_{i} x_{i}^{2} = \frac{1}{2} w_{i} x_{i} (l_{i} - x_{i})$$
 . . . (115)

For a uniform load over all, the moment at the center of the girder becomes,

$$M_c = \frac{1}{2} w_i \frac{l_i^2}{2} - \frac{1}{2} w_i \frac{l_i^2}{4} = \frac{1}{8} w_i l_i^2 \dots \dots (116)$$

The well known formula for this case. We see, therefore, that our formulas are perfectly exact for the simple girder, and also that a variable moment of inertia or difference of level of the supports does not effect the values of the bending moments or the shears.

From (E), we have,

$$y_{i} = h_{i} + t_{i} x_{i} + \frac{1}{6 E I_{x}^{i}} \left(S_{i} x_{i}^{3} - \Sigma P_{i} (x_{i} - a_{i})^{3} \right)$$

$$+ \frac{1}{6 E I_{x}^{i}} \sum_{v = x_{i}} \left(\frac{I_{x}^{i}}{I_{v-i}^{i}} - \frac{I_{x}^{i}}{I_{v}} \right) \left\{ S_{i} e_{v}^{2} (3 x_{i} - 2 e_{v}) - \Sigma P_{i} (e_{v} - a_{i})^{3} - 3 (x_{i} - e_{v}) \Sigma P_{i} (e_{v} - a_{i})^{2} \right\}$$

$$(117)$$

(b) E and I constant.

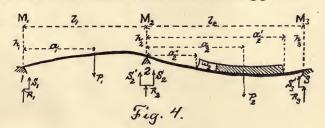
The moment and shear equations are the same as in (a).

$$y_i = h_i + t_i x_i + \frac{1}{6 E I} \left\{ S_i x_i^3 - \Sigma P_i (x_i - a_i)^3 \right\} \dots \dots (119)$$

If the origin is taken at one of the supports, the corresponding value of h, will, of course, equal zero.

CASE II.

A beam continuous over three supports'-



(a) E alone, constant.

$$M_{2} = \frac{-c_{2}}{d_{3} \beta_{i}^{"}} \left\{ (A_{2} + Y_{2} + X_{2}^{'} + X_{i}^{"}) d_{2} + B_{2} d_{i} \right\} + \frac{-d_{2}}{c_{3} \beta_{2}} \left\{ (A_{i} + Y_{i} + X_{i}^{'}) c_{i} + B_{i} c_{2} \right\} . . (121)$$

Substituting (122) and (123) in (121), it reduces to, $\frac{1}{2}(A+V+V'+V''+P)$

 $Y_{2} = -36 E^{2} \mathring{I}_{t} \mathring{I}_{t} \left\{ \frac{h_{2} - h_{1}}{L} + \frac{h_{2} - h_{3}}{L} \right\} \qquad (134)$

A2 and B1

Concentrated loads-

$$A_{2} = - \Sigma P_{2} l_{2}^{2} (2 k_{2} - 3 k_{2}^{2} + k_{2}^{3}) 6 E I_{1}^{'} \dots \dots \dots \dots (135)$$

$$B_{i} = - \Sigma P_{i} l_{i}^{z} (k_{i} - k_{i}^{s}) 6 E \hat{l}_{i}^{z} \dots \dots \dots \dots (136)$$

Partial uniform loads-

$$A_{2} = -w_{2} l_{2}^{3} \left\{ k_{2}^{2} - k_{2}^{2} + \frac{k_{2}^{4}}{4} \right\} \begin{cases} 6 E I_{1} \\ a_{2}^{\prime} = k_{2} l_{2} \end{cases}$$

$$(137)$$

$$B_{i} = -w_{i} l_{i}^{s} \left\{ \frac{2 k_{i}^{s} - k_{i}^{s}}{4} \right\} a_{i}^{s} = k_{i} l_{i} (138)$$

$$a_{i}^{s} = k_{i} l_{i}$$

Uniform load over all-

Shears and intermediate bending moments-

For shears and the intermediate bending moments, apply the general equations (B) (c) and (D), and for deflection use equations (E) and (45).

(b) E and I constant.

In this case, we have from the equations under (a),

$$M_2 = \frac{Y_2 + A_2 + B_1}{2(l_1 + l_2)}$$
 (142)

In which,

$$Y_2 = -6 E I \left\{ \frac{h_2 - h_1}{l_1} + \frac{h_2 - h_3}{l_2} \right\} \dots \dots (143)$$

A2 and B1

Concentrated loads-

Uniform load over all-

$$A_{2} = -w_{2} l_{2}^{3} \left\{ k_{2}^{2} - k_{2}^{3} + \frac{k_{2}^{4}}{4} \right\} \begin{matrix} a'_{2} = k_{2} l_{2} \\ \vdots \\ a''_{2} = k_{2} l_{2} \end{matrix} . \dots (146)$$

$$B_{i} = -w_{i} l_{i}^{s} \left\{ \frac{2 k_{i}^{s} - k_{i}^{s}}{4} \right\} \frac{a_{i}^{s} = k_{i} l_{i}}{a_{i}^{s} = k_{i} l_{i}} (147)$$

Uniform load over all-

As this is a case much used in practice, we will give expressions for the bending moments in terms of the spans and the loads.

Concentrated loads-

$$M_{2} = \frac{Y_{2} - \sum P_{2} l_{2}^{2} K_{2} - \sum P_{i} l_{i}^{2} K'_{i}}{2 (l_{i} + l_{0})} (150)$$

In which,

And

From
$$(B)$$
 and (C) ,
 $Y_{\circ} - \stackrel{\Sigma}{\Sigma} P_{\circ} L^{\circ} K_{\circ} - \stackrel{\Sigma}{\Sigma} P_{\circ} L K'$

$$S_{i} = \frac{Y_{2} - \Sigma P_{2} l_{2}^{2} K_{2} - \Sigma P_{i} l_{i} K_{i}^{\prime}}{2 l_{i} (l_{i} + l_{2})} + \Sigma P_{i} (1 - k_{i}) = R_{i}. \quad (153)$$

$$S_{z}' = \frac{-Y_{z} + \sum P_{z} l_{z}^{2} K_{z} + \sum P_{i} l_{i}^{2} K_{i}}{2 l_{i} (l_{i} + l_{z})} + \sum P_{i} k_{i} ...$$
(154)

$$S_{z}' = \frac{-Y_{z} + \sum P_{z} l_{z}^{2} K_{z} + \sum P_{i} l_{i}^{2} K_{i}}{2 l_{i} (l_{i} + l_{z})} + \sum P_{i} k_{i} . .$$

$$S_{z} = \frac{-Y_{z} + \sum P_{z} l_{z}^{2} K_{z} + \sum P_{i} l_{i}^{2} K_{i}}{2 l_{z} (l_{i} + l_{z})} + \sum P_{z} (1 - k_{z})$$

$$= R_{z}$$

$$(154)$$

Or,

$$R_{2} = \frac{-Y_{2} + \Sigma P_{2} l_{2}^{2} K_{2} + \Sigma P_{i} l_{i}^{2} K_{i}}{2 l_{i} l_{2}} + \Sigma P_{i} l_{i}^{2} K_{i}^{2} + \Sigma P_{i} l_{i} + \Sigma P_{2} (1 - k_{2}) (156)$$

$$S_{s}' = \frac{Y_{z} - \sum P_{z} l_{z}^{2} K_{z} - \sum P_{t} l_{t}^{2} K_{t}}{2 l_{z} (l_{t} + l_{z})} + \sum P_{z} k_{z} = R_{z} (157)$$

Uniform loads over each span-

The above equations at once reduce to

$$M_{2} = \frac{Y_{2} - \frac{1}{4} w_{2} l_{2}^{3} - \frac{1}{4} w_{1} l_{1}^{3}}{2(l_{1} + l_{2})} \dots \dots \dots \dots \dots (158)$$

$$S_{i} = \frac{Y_{2} - \frac{1}{4} w_{2} \ l_{2}^{3} - \frac{1}{4} w_{i} \ l_{i}^{3}}{2 \ l_{i} \ (l_{i} + l_{2})} + \frac{1}{2} w_{i} \ l_{i} = R_{i} \quad . \quad . \quad (159)$$

$$S_{2}' = \frac{-Y_{2} + \frac{1}{4} w_{2} l_{2}^{3} + \frac{1}{4} w_{i} l_{i}^{3}}{2 l_{i} (l_{i} + l_{2})} + \frac{1}{2} w_{i} l_{i}. \quad (160)$$

$$S_{2} = \frac{-Y_{2} + \frac{1}{4} w_{2} l_{2}^{3} + \frac{1}{4} w_{1} l_{1}^{3}}{2 l_{2} (l_{1} + l_{2})} + \frac{1}{2} w_{2} l_{2} (161)$$

Or.

$$R_{2} = \frac{-Y_{2} + \frac{1}{4} w_{2} l_{2}^{3} + \frac{1}{4} w_{i} l_{i}^{3}}{2 l_{i} l_{2}} + \frac{1}{2} (w_{i} l_{i} + w_{2} l_{2}) . . . (162)$$

$$S_{3}' = \frac{Y_{2} - \frac{1}{4} w_{2} l_{3}^{3} - \frac{1}{4} w_{i} l_{i}^{3}}{2 l_{2} (l_{i} + l_{2})} + \frac{1}{2} w_{2} l_{2} - R_{3} \dots$$
 (163)

(c) E, I and h constant.

In this case, the formulas are the same as in case (b) with the term $Y_s = 0$.

(d) E, I, h and I constant.

Making $l_1 = l_2$, the equations of case (b) reduce to

Concentrated loads-

$$M_2 = \frac{-\Sigma P_2 l_2 K_2 - \Sigma P_i l_i K_i'}{4} \qquad (164)$$

$$S_{i} = \frac{-\Sigma P_{s} K_{s} - \Sigma P_{i} K_{i}}{4} + \Sigma P_{i} (1 - k_{i}) = R_{i} \dots \dots (165)$$

$$S_{s}' = \frac{+\Sigma P_{s}}{4} \frac{K_{s} + \Sigma P_{s} K_{s}'}{4} + \Sigma P_{s} k_{s} ...$$
(166)

$$S_{2}' = \frac{+ \Sigma P_{2} K_{2} + \Sigma P_{i} K_{i}'}{4} + \Sigma P_{i} k_{i} ...$$

$$S_{2} = \frac{+ \Sigma P_{2} K_{2} + \Sigma P_{i} K_{i}'}{4} + \Sigma P_{2} (1 - k_{2})$$

$$(166)$$

$$R_{s} = \frac{\sum P_{s} K_{s} + \sum P_{t} K'_{t}}{2} + \sum P_{t} k_{t} + \sum P_{s} (1 - k_{s}) = R_{s} \quad . \quad . \quad (168)$$

$$S_s' = \frac{-\Sigma P_s K_s - \Sigma P_t K_t'}{4} + \Sigma P_s k_s = R_s \dots \dots (169)$$

Uniform load over all-w,-w,-w

$$S_{l} = -\frac{1}{8} w l + \frac{1}{2} w l = \frac{3}{8} w l = R_{l} \dots$$
 (171)

$$S_{s} = + \frac{1}{8} w l + \frac{1}{2} w l = \frac{5}{8} w l$$

$$S_{s} = + \frac{1}{8} w l + \frac{1}{2} w l = \frac{5}{8} w l$$

$$(172)$$

$$= \frac{10}{8} w l = R_{s}$$

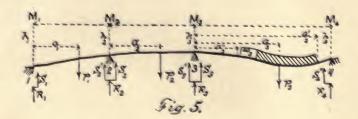
$$\dots \dots (173)$$

$$S_2 = +\frac{1}{8} w l + \frac{1}{2} w l = \frac{5}{8} w l$$
 (173)

$$S_s = -\frac{1}{8}wl + \frac{1}{2}wl = \frac{3}{8}wl = R_s \dots (174)$$

CASE III.

A beam continuous over four supports-



(a) E alone, constant.

$$M_{e} = \frac{-c_{e}}{d_{s} \, \tilde{\beta}_{i}} \left\{ \frac{(A_{e} + Y_{e} + X_{e} + X_{i}^{*}) \, d_{s} + B_{e} \, d_{e}}{(A_{s} + Y_{s} + X_{s} + X_{e}^{*}) \, d_{e} + B_{s} \, J_{e}} + \frac{-d_{s} \, \tilde{\beta}_{e}^{*}}{c_{s} \, \tilde{\beta}_{s} \, \tilde{\beta}_{e}} \left\{ (A_{i} + Y_{i} + X_{i}^{*} + X_{e}^{*}) \, c_{i} + B_{i} \, c_{e} \right\} . \quad (176)$$

Which reduces to the following:

$$M_{z} = \frac{-I}{d_{z} \, \delta_{r}} \left\{ \frac{\left(A_{z} + Y_{z} + X_{z} + X_{r}^{*}\right) \, d_{z} + B_{z}}{\left(A_{z} + Y_{z} + X_{z}^{*} + X_{z}^{*}\right)} \right\} + \frac{-d_{z}}{\left(c_{z} \, \delta_{z}\right)} \frac{\delta_{z}^{*}}{\delta_{z}} \, B_{r} \dots (177)$$

$$M_{s} = \frac{-c_{s} \beta_{z}}{(d_{s} \overline{\beta_{r}}) \beta_{z}^{2}} \left\{ (A_{s} + Y_{s} + X_{s}^{\prime} + X_{z}^{\prime}) \right\} + \frac{-1}{c_{s} \beta_{s}} \left\{ (A_{z} + Y_{z} + X_{z}^{\prime} + X_{r}^{\prime}) + B_{z} c_{s} \right\} . . . (178)$$

Use general equations for the various terms in the above, and also for shear, intermediate moments and deflection.

(b) **E** and **I** constant.

Our equations now reduce to

$$M_{z} = \frac{(Y_{z} + A_{z} + B_{1}) 2 (l_{z} + l_{3}) - (A_{3} + B_{z} + Y_{3}) l_{z}}{4 (l_{1} + l_{2}) (l_{z} + l_{3}) - l_{z}^{2}} (179)$$

$$M_{3} = \frac{(Y_{3} + A_{3} + B_{2}) 2 (l_{1} + l_{2}) - (B_{1} + A_{2} + Y_{2}) l_{2}}{4 (l_{1} + l_{2}) (l_{2} + l_{3}) - l_{2}^{2}} . . . (180)$$

(c) E, I and h constant.

$$M_{2} = \frac{2(l_{2}+l_{3})}{4} \frac{(A_{2}+B_{1})-l_{2}(A_{3}+B_{2})}{(l_{i}+l_{2})(l_{2}+l_{3})-l_{2}^{2}} \dots \dots (181)$$

$$M_{3} = \frac{2 (l_{1}+l_{2}) (A_{3}+B_{2})-l_{2} (A_{2}+B_{1})}{4 (l_{1}+l_{2}) (l_{2}+l_{3})-l_{2}^{2}} \dots \dots (182)$$

(d) E, I, h and l constant.

$$M_s = \frac{4 (A_s + B_z) - (A_s + B_s)}{15 l} \dots \dots \dots (184)$$

Uniform loads-

If each span is covered with a uniform load, we have

If $w_1 = w_2 = w_3 = w$, and $l_1 = l_2 = l_3 = l$, we obtain, at once, the well known forms,

And for the shears we have

$$S_{t} = -\frac{1}{10} w l + \frac{1}{2} w l = \frac{4}{10} w l = R_{t} \dots (189)$$

$$S_2 = + \frac{1}{10} w l + \frac{1}{2} w l = \frac{6}{10} w l \dots (190)$$

$$S_2 = +\frac{1}{2} w l$$
 = $\frac{5}{10} w l \dots (191)$

$$R_2 = S_1' + S_2 = \left\{ \frac{6}{10} + \frac{5}{10} \right\} w l = \frac{11}{10} w l \dots (192)$$

$$R_s = S_s' + S_s = \left\{ \frac{5}{10} + \frac{6}{10} \right\} w l = \frac{11}{10} w l \dots (195)$$

$$S_{l} = -\frac{1}{10} w l + \frac{1}{2} w l = \frac{4}{10} w l = R_{l} \dots (196)$$

CASE IV.

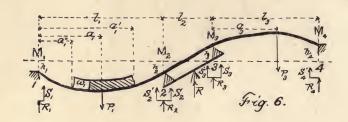
"THE TIPPER."

* We will now consider the case of the "Tipper," or a beam continuous over four supports, the second and third supports, resting on a rigid beam, supported, generally, by a single support at its center.

It is evident that if supports 2 and 3 are supported by an unyielding bar, which is supported, in turn, at its center, that the reaction R_2 must equal the reaction R_3 , and also that $+h_2=-h_3$. If the unyielding bar is not supported at its center, R_2 and R_3 will be inversely proportional to their lever arms around the point of support, and $+h_2$ and $-h_3$ will be directly proportional to the lever arms.

^{*} See "Strains in Framed Structures," page 254, by Prof. DuBois.

The unyielding bar supported at its center-



(a) E alone, constant.

From our general equations, by reduction, we obtain,

$$R_{2} = \frac{-M_{2}}{l_{i}} + \frac{M_{3} - M_{2}}{l_{2}} + Q_{i} + Q_{2} + Q_{3} + Q_{4} + Q_{5} +$$

Since $R_s = R_s$, we have

$$\frac{M_{s}}{l_{s}} - \frac{M_{z}}{l_{t}} + \frac{2 M_{s} - 2 M_{z}}{l_{z}} + Q_{t} + Q_{z} - Q_{z} - Q_{z} = 0 \quad . \quad . \quad (199)$$

Let

Then (199) becomes

$$n M_s - m M_z + 2 m n (M_s - M_z) + m n Q l_z = 0$$
 (203)
Or,

$$n M_3 - m M_2 + 2 m n (M_3 - M_2) = -m n Q l_2 \dots (204)$$

By inspecting the equations under Case III, we see that the moment equations may be placed under the following forms:

From (205) and (206), we have

$$M_3 - M_2 = (C_2' - C_2) Y_2 + (C_3' - C_3) Y_3 + (C_1' - C_1) \dots (207)$$

Substituting (205), (206) and (207) in (204), it becomes

$$\begin{cases}
+n C_{2}' Y_{2}+n C_{3}' Y_{3}+n C_{i}' \\
-m C_{2} Y_{2}-m C_{3} Y_{3}-m C_{i} \\
+2 m n (C_{3}'-C_{2}) Y_{2}+2 m n (C_{3}'-C_{3}) Y_{3} \\
+2 m n (C_{i}'-C_{i})
\end{cases} = -m n Q l_{2}. (208)$$

Or, by placing V and V' in place of the co-efficients of Y, etc., we have

Now,
$$Y_2 = -\theta_1 \theta_2 \left\{ \frac{h_2 - h_1}{n l_2} + \frac{h_2 - h_3}{l_2} \right\} \dots \dots (210)$$

$$Y_{s} = -\theta_{s} \, \theta_{s} \left\{ \frac{h_{s} - h_{s}}{l_{s}} + \frac{h_{s} - h_{s}}{m \, l_{s}} \right\} \quad . \quad . \quad . \quad . \quad . \quad . \quad (211)$$

But, $h_2 = -h_3$, or, $h_3 = -h_2$; therefore,

$$V Y_{2} = -\theta_{i} \theta_{2} \left\{ \frac{-h_{i}}{n l_{2}} V \right\} - \left\{ \frac{1+2n}{n l_{2}} V \right\} \theta_{i} \theta_{2} h_{2} ... (212)$$

$$V' Y_3 = -\theta_2 \theta_3 \left\{ \frac{-h_4}{m l_2} V' \right\} - \left\{ \frac{1+2m}{m l_2} V' \right\} \theta_2 \theta_3 h_2 . . (213)$$

Hence,

$$V' Y_{s} + V Y_{z} = \left[\left\{ \frac{h_{s}}{n l_{z}} V \theta_{s} \right\} + \left\{ \frac{h_{s}}{m l_{z}} V' \theta_{s} \right\} \right] \theta_{z}$$

$$- \left[\left\{ \frac{1 + 2 n}{n l_{z}} V \theta_{s} \right\} + \left\{ \frac{1 + 2 m}{m l_{z}} V' \theta_{s} \right\} \right] \theta_{z} h_{z} . . (214)$$

Therefore,

[1]
$$\theta_2$$
—[2] θ_2 h_2 = — $m \ n \ Q \ l_2$ — V'' (216)

And

$$h_2 = \frac{+m \ n \ Q \ l_2 + V'' + [1] \ \theta_2}{[2] \ \theta_2} \dots \dots \dots (217)$$

From (217), we can find the value of $h_2 = -h_3$, and then the values of Y_2 and Y_3 from (210) and (211), whence we can determine M_2 and M_3 from (177) and (178).

The easiest method of finding the values of the terms in the equation giving the value of h_2 , is to substitute, for any particular case, the values of all the known quantities in the immediately preceding equations, or find the values of the constants C_i , C_2 , etc., and in turn, substitute them in the equations containing them. The operations are long and exceedingly tedious, but simple, as an examination of the equations shows.

(b) **E** and **I** constant.

If the moment of inertia is constant, the deduction of the constants C_i , C_2 , etc., becomes much more simple. After h_2 is determined from (217), and Y_2 and Y_3 from (210) and (211), the bending moments are readily obtained from (179) and (180).

If $h_i = h_a = 0$, as is usually the case, the process becomes still more simple.

(c) **E** and **I** constant.
$$h_i = h_i = 0$$
. $l_i = l_3 = l$. $l_g = n \ l$. $h_g = -h_g$. We have, as in (a), by a little reduction, $n \ (M_s - M_g) + 2 \ (M_g - M_g) = -n \ l \ Q$ (218) Or, $(n+2) \ (M_g - M_g) = -n \ l \ Q$ (219) Letting $4 \ (l_i + l_g) \ (l_2 + l_g) - l_g^2 = D \ l$ (220) We have, from (179) and (180), $M_g = \frac{(Y_g + A_g + B_i) \ 2 \ (n+1) - (A_g + B_g + Y_g) \ n}{D}$ (221) $M_g = \frac{(Y_g + A_g + B_g) \ 2 \ (n+1) - (A_g + B_i + Y_g) \ n}{D}$ (222)

Then, $D(M_3-M_2)=(Y_3+A_3+B_2) \ 2(n+I)-(A_2+B_1+Y_2) \ n$ $(Y_3+A_3+B_2) \ n -(A_2+B_1+Y_2) \ 2(n+I) \ . \ . \ (223)$ Hence,

$$Y_s (3 n+2)-Y_s (3 n+2)+(A_s+B_s -A_s-B_s) (3 n+2)=(M_s-M_s) D$$
. (224)

Substituting (224) in (219), it reduces to

$$(Y_3 - Y_2 + A_3 + B_2 - A_2 - B_1)(3n+2)(n+2) = -DnlQ$$
. (225)

Therefore,

$$Y_s - Y_z = \frac{-D \ n \ l \ Q}{(3 \ n + 2) \ (n + 2)} - A_s - B_z + A_z + B_t \quad . \quad . \quad (226)$$

Since $h_1 = h_2 = 0$ and $h_2 = -h_3$,

$$Y_{z} = -6 E I \left\{ \frac{h_{z}}{l_{z}} + \frac{2 h_{z}}{l} \right\} = -6 E I \left\{ \frac{h_{z}(1+2 n)}{n l} \right\}.$$
 (227)

$$Y_{3} = -6 E I \left\{ \frac{-h_{2}}{l} + \frac{-2 h_{2}}{l_{2}} \right\} = +6 E I \left\{ \frac{h_{2}(1+2n)}{n l} \right\} (228)$$

Therefore,

And

Hence,

$$Y_{2} = \frac{D \ n \ l \ Q}{2 \ (3 \ n+2) \ (n+2)} + \frac{A_{3} + B_{2} - A_{2} - B_{1}}{2} \ \dots \ (231)$$

But,
$$D \stackrel{!}{l=4} (l_1+l_2) (l_2+l_3)-l_2^2=4 l$$
, $l_3+4 l_2 l_3$
 $+4 l_2 l_3+4 l_2^2-l_2^2=4 l^2+8 n l^2+3 n^2 l^2=l^2 (3 n^2 +8 n+4)=l^2 (3 n+2) (n+2) \dots (232)$

And

Substituting (201) and (232) in (231), we have

$$Y_{z} = \frac{l l_{z}(Q'_{z} + Q_{z} - Q'_{z} - Q_{s}) + A_{z} + B_{z} - A_{z} - B_{t}}{2} = -Y_{s} . . (233)$$

(233) completely determines Y_z and Y_s , and now the bending moments can be deduced from (179) and (180).

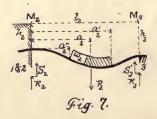
III.

BEAMS WITH FIXED ENDS.

In this chapter we shall consider beams with one or both ends fixed.

CASE I.

A beam fixed at one end and supported at the other—



For convenience, the left end will be considered as fixed.

This case is, in reality, the same as Case II, in Chapter II, considering $l_t=0$, $h_t=h_z$ and $f_t=0$, hence, we can use the equations of that case by making the proper changes.

(a) **E** alone constant.

(124) becomes

(134) becomes

By inspecting equations (124) to (134), inclusive, it is seen that we can cancel out $\mathbf{6} \mathbf{E} \mathbf{I}_{i}^{'}$ from all the terms in (234), hence, we can write

$$v=1$$

$$F'_{2} = \sum_{v=l_{2}} \stackrel{2}{\bigtriangleup}_{v} e_{r}^{2} \left\{ \frac{3}{l_{2}} - \frac{2}{l_{2}^{2}} \right\} \qquad (131)$$

$$H_{2} = \sum_{v=l_{2}}^{2} \triangle_{v}^{2} \left\{ \frac{\sum_{v} P_{2} (e_{v} - a_{2})^{3}}{l_{2}} + \frac{3(l_{2} - e_{v})}{l_{2}} \sum_{v} P_{2} (e_{v} - a_{2})^{2} \right\} ...(127)$$

Concentrated loads-

$$A_{z} = - \Sigma P_{z} l_{z}^{z} (2 k_{z} - 3 k_{z}^{z} + k_{z}^{s}) \dots \dots (14 + 1)$$

^{*} In reality, $I_i\left\{\frac{h_i-h_2}{l_i}\right\} = I_i^{j}\left\{\frac{\theta_i}{\theta_2}\right\}$ which is indeterminate, but the above form is the only logical one which the expression for X_2 can take.

Partial uniform loads-

$$A_{2} = -w_{2} l_{2}^{3} \left\{ k_{2}^{2} - k_{2}^{3} + \frac{k_{2}^{4}}{4} \right\} \begin{matrix} a'_{2} = k_{2} l_{2} \\ \vdots \\ a''_{n} = k_{n} l_{n} \end{matrix} . \dots (146)$$

Uniform load over all-

Shears and intermediate bending moments-

For shears and the intermediate bending moments, apply the general equations (B), (C) and (D), and for deflection, use equations (E) and (45).

(b) E and I constant.

For this case, (129), (131), (133), (127) and (238) become zero, and (237) reduces to

And we have from (234) or (142)

In which the values of Y_z and A_z are given by (236), (144), (146) or (148).

As this is a case quite likely to occur in practice, we will give expressions for the bending moments and shears in terms of the loads and span.

Concentrated loads-

Substituting in (144) in (240), or from (150), we have

$$M_2 = \frac{Y_2 - \Sigma P_2 l_2^2 K_2}{2 l_0} \dots (241)$$

In which $K_2 = 2 k_2 - 3 k_2^2 + k_2^3$.

From (B) and (C), or (155),

$$S_{2} = \frac{-Y_{2} + \sum P_{2} l_{2}^{2} K_{2}}{2 l_{2}^{2}} + \sum P_{2} (1 - k_{2}) = R_{2} (242)$$

Also, (see 157).

$$S_{s} = \frac{Y_{2} - \sum P_{2} l_{2}^{2} K_{2}}{2 l_{s}^{2}} + \sum P_{2} k_{2} = R_{s} \dots (243)$$

Uniform load over all-

The above equations at once reduce to

$$S_{2} = \frac{-Y_{2} + \frac{1}{4} w_{2} l_{2}^{3}}{2 l_{2}^{2}} + \frac{1}{2} w_{2} l_{2} = R_{2} \dots \dots \dots (245)$$

$$S_{3}^{\prime} = \frac{+ Y_{2} - \frac{1}{4} w_{2} l_{2}^{3}}{2 l_{2}^{2} + \frac{1}{2} w_{2} l_{2} = R_{2} \dots (246)}$$

(c) E, I and h constant.

In this case the formulas are the same as in case (b) with the term $Y_2=0$.

The formulas are:

Concentrated loads-

$$S_2 = \frac{\sum P_2 K_2}{2} + P_2 (1-k_2) = R_2 \dots \dots (248)$$

$$S_{s} = \frac{-\Sigma P_{2} K_{2}}{2} + \Sigma P_{2} k_{s} = R_{s} \dots$$
 (249)

Uniform load over all-

(244) becomes

$$S_z = \frac{1}{8} w_z l_z + \frac{1}{2} w_z l_z = \frac{5}{8} w_z l_z = R_z \dots \dots (251)$$

$$S_{3}' = -\frac{1}{8} w_{2} l_{2} + \frac{1}{2} w_{2} l_{2} = \frac{3}{8} w_{2} l_{2} = R_{3} \dots (252)$$

From (18), we have

From (E_t) and (45 a), we obtain, if $h_2 = h_3 = 0$,

$$y_z = -\frac{1}{48 E I} w_z x_z^2 (3 l_z^2 - 5 l_z x_z + 2 x_z^2) ... x_z = a_z'... (255)$$

Special case-

A single concentrated load at the centre of the beam.

By applying Table I, (250) becomes at once

$$M_2 = -P_2 l_2 \ 0.1875 = -\frac{3}{16} P_2 l_2 \dots (256)$$

$$S_2 = \frac{3}{16} P_2 + \frac{1}{2} P_2 = \frac{11}{16} P_2 = R_2 \dots (257)$$

$$S_s = -\frac{3}{16} P_z + \frac{1}{2} P_z = \frac{5}{16} P_z = R_s \dots$$
 (258)

$$M_{x}^{2} = -\frac{3}{16} P_{z} l_{z} + \frac{11}{16} P_{z} x_{z} - P_{z} \left\{ x_{z} - \frac{1}{2} l_{z} \right\} x_{z} = a_{z} = \frac{1}{2} l_{z}$$

Or,

$$M_x^2 = \frac{5}{16} P_2 (l_2 - x_2) \dots x_2 > \frac{1}{2} l_2 \dots (259)$$

$$M_x^2 = \frac{1}{16} P_2 (11 x_2 - 3 l_2) \dots x_2 < \frac{1}{2} l_2 \dots (260)$$

The maximum moment occurs when $x_2 = \frac{1}{2} l_2$, or,

Max.
$$M_x^2 = \frac{5}{32} P_2 l_2 \dots (261)$$

From
$$(\boldsymbol{E}_{1})$$
, if $h_{2}=h_{3}=0$,

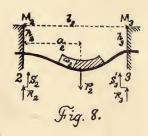
$$y_{z} = \frac{1}{96 E I} P_{z} (15 l_{z} x_{z}^{2} - 5 x_{z}^{2} - 12 l_{z}^{2} x_{z} + 2 l_{z}^{3}) . . x_{z} > \frac{1}{2} l_{z} . . (262)$$

And

$$y_2 = \frac{1}{96 E I} P_2 x_2^2 (11 x_2 - 9 l_2) \dots x_2 < \frac{1}{2} l \dots (263)$$

CASE II.

A beam fixed at both ends-



This is the same as Case II, Chapter II, with $h_i = h_2$, $h_4 = h_3$, $l_i = 0$, $l_3 = 0$, and $I_i^j = 0$.

(a) E alone constant.

From (177), we have

$$M_{2} = \frac{1}{d_{A} \beta'_{i}} \left\{ \frac{(A_{2} + Y_{2} + X_{2}') d_{3} + B_{2}}{(+Y_{3} + X_{2}'')} \right\} \dots \dots (264)$$

From (178),

$$M_{s} = \frac{-c_{s} \beta_{2}}{(d_{s} \beta_{s}^{"}) \beta_{2}^{"}} \left\{ +Y_{s} + X_{s}^{"} \right\} + \frac{-1}{c_{s} \beta_{s}} \left\{ (A_{2} + Y_{2} + X_{s}^{'}) + B_{s} c_{s} \right\} . . (265)$$

(264) reduces to

$$M_{2} = \frac{-\beta_{2}''}{4\beta_{2}'\beta_{3}' - \beta_{2}'\beta_{2}''} \left\{ \frac{(A_{2} + Y_{2} + X_{2}') - \beta_{2}''}{(B_{2} + Y_{3} + X_{2}'')} \right\} ... (266)$$

And (178) becomes

$$M_{s} = \frac{2 \beta_{2}'}{4 \beta_{2}' \beta_{3}' - \beta_{2} \beta_{2}''} \left\{ Y_{s} + X_{2}'' \right\} + \frac{-\beta_{2}}{4 \beta_{2}' \beta_{3}' - \beta_{2} \beta_{2}''} \left\{ (A_{2} + Y_{2} + X_{2}') + (B_{2}) \frac{-2 \beta_{2}'}{\beta_{2}} \right\} . (267)$$

In which, after dividing by $\theta_1 = \theta = \theta_s$,

$$\beta_3' = l_2 + F_2''$$
 . . . The values of F can be found from (269)

The values of the remaining terms are easily found from the general equations, remembering that θ_1 or θ_2 has been cancelled out of each term.

(b) E and I constant.

From (266) or (179), we have

$$M_{2} = \frac{(Y_{z} + A_{z}) 2 l_{z} - (B_{z} + Y_{z}) l_{z}}{3 l_{z}^{2}} = \frac{2 (Y_{z} + A_{z}) - (B_{z} + Y_{z})}{3 l_{z}} . . (274)$$

From (267) or (180), we obtain

$$M_{s} = \frac{(Y_{s} + B_{2}) 2 l_{2} - (A_{z} + Y_{2}) l_{2}}{3 l_{z}^{2}} = \frac{2 (Y_{s} + B_{2}) - (A_{z} + Y_{2})}{3 l_{s}} . . (275)$$

In which the terms have the values given under (a).

^{*} See note under Equation (235).

(c) E, I and h constant.

(274) becomes

(275) becomes

For concentrated loads we can write

$$M_{z} = \frac{-2 \Sigma P_{z} l_{z} K_{z} + \Sigma P_{z} l_{z} K_{z}'}{3} \dots \dots \dots (278)$$

And

Uniform load over all-

$$M_x^2 = w_2 \left\{ \frac{l_2 x_2}{2} - \frac{x_2^2}{2} - \frac{l_2^2}{12} \right\} \dots \dots \dots (282)$$

If x=0

Max.
$$M_x^2 = -\frac{1}{12} w_2 l_2^2 = M_2 = M_3 \dots \dots \dots \dots (283)$$

If $h_s = h_s = 0$,

$$y_{z} = \frac{-1}{24 E I} w_{z} x_{z}^{2} (l_{z}^{2} - 2 l_{z} x_{z} + x_{z}^{2}) \dots x_{z} = a'_{z} \dots (284)$$

A single load in the centre of the beam-

By using Table I, we have at once from (278) and (279),

$$M_z = M_s = -0.125 P_z l_z = -\frac{1}{8} P_z l_z \dots \dots (285)$$

$$M_x^2 = \frac{1}{8} P_2 (3 l_2 - 4 x_2) \dots x_2 > a_2 = \frac{1}{2} l_2 \dots (287)$$

Also,

$$M_x^2 = \frac{1}{8} P_2 (4 x_2 - l_2) \dots x_2 < \frac{1}{2} l_2 \dots (288)$$

If
$$x_2=0$$
, or $\frac{1}{2}$ l_2 .

Max.
$$M_x = -\frac{1}{8} w_2 l_2$$
, or $+\frac{1}{8} w_2 l_2$ (289)

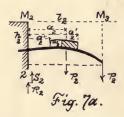
$$y_2 = \frac{-1}{48 E I} P_2 (4 x_2^3 + 6 l_2^2 x_2 - 9 l_2 x_2^2 - l_2^3) . . x_2 > a_2 = \frac{1}{2} l_2 . (290)$$

And

$$y_2 = \frac{-1}{48EI} P_2 x_2^2 (3 l_2 - 4 x_2) ... x_2 < a_2 = \frac{1}{2} l_2 ... (291)$$

CASE III.

A beam fixed at one end, and unsupported at the other—



As the right end of the beam is unsupported, there can be no reaction $S_s = R_s$, therefore, by (11),

Concentrated loads-

$$S_s'=rac{M_z-M_s}{l_z}+\Sigma\,P_z\,k_z{=}0,$$
 and since $M_s{=}0,$ we have

$$M_2 = -\Sigma P_2 l_2 k_2 \dots S_2 = \Sigma P_2 \dots \dots (292)$$

Showing that the moment of inertia does not enter into the expression for the bending moment.

Any uniform load-

$$M_z = -\frac{1}{2} w_z (a'_z - a''_z) (a'_z + a''_z) \dots \dots \dots (293)$$

From (8),

$$M_x = 0.$$
 $x_2 = a_2$

And

$$M_x^2 = \Sigma P_2 (x_2 - a_2) \dots x_2 < a_2 \dots \dots (295)$$

Also, for any uniform load,

$$M_x^2 = w_2 \left(a_2' - a_2'' \right) \left\{ x_2 - \frac{a_2' + a_2''}{2} \right\} \dots x_2 \leq a_2'' \dots (296)$$

For deflection, use the general formulas, if the moment of inertia is variable.

If the moment of inertia is constant, we have, from (\mathbf{E}_i) , If $h_2=0$,

$$y_2 = + \frac{1}{6 E I} \left\{ 3 M_2 x_2^2 + S_2 x_2^3 - \Sigma P_2 (x_2 - a_2)^3 \right\} (297)$$

Or, for a single concentrated load,

$$y_{2} = \frac{P_{2}}{6 E I} \left\{ -\beta a_{2} x_{2}^{2} + x_{2}^{3} - (x_{2} - a_{2})^{3} \right\} \quad . \quad . \quad . \quad . \quad (298)$$

And if the load is at the end of the beam, $x_2=a_2=l_2$, and (298) becomes

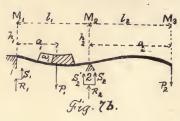
$$y_2 = -\frac{P_2 l_2^3}{3 E I}$$
...at end of beam....(299)

Also, for a uniform load over all, $(x_2=a_2=l_2)$,

$$y_z = \frac{-1}{8EI} w_z l_z^A \dots \text{ at end of beam } \dots \dots \dots (300)$$

CASE IV.

A beam on two supports, and one end unsupported-



In this case, M_1 , M_3 and S_3 equal zero, hence, from (11), or (C), we have

$$S_s' = \frac{M_z}{l_z} + Q_s' = 0$$
, or, $M_z = -Q_s' l_z \dots \dots \dots (301)$

Which is precisely the same equation we obtained in Case The values of Q'_{2} are given in (95), (97) and (99).

From the general equations (B) and (C), we obtain

$$S_{z}' = \frac{-M_{z}}{l_{s}} + Q_{s}' \qquad (303)$$

$$S_{2} = \frac{-M_{2}}{l_{1}} + Q_{1}'$$

$$= R_{2}$$

$$S_{2} = \frac{-M_{2}}{l_{2}} + Q_{2}$$

$$(303)$$

From (D), we obtain

$$M_x = S_i x_i - L_i = \frac{M_2 x_i}{L} + Q_i x_i - L_i \dots (305)$$

$$M_{x}^{2} = M_{z} + S_{z} x_{z} - L_{z} = M_{z} + \frac{-M_{z} x_{z}}{l_{z}} + Q_{z} x_{z} - L_{z} \dots (306)$$

Note that the moment of inertia does not appear in any of the above equations.

For deflection, use the general equation (E).

CASE V.

A beam on one support, having one end fixed, and the other unsupported—

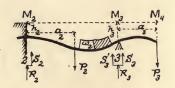


Fig. 8a.

(a) E.alone, constant.

The shears, intermediate bending moments and deflections are readily obtained from the general equations.

. (307)

 $M_3 = -Q_3' l_3 \dots$

CASE VI.

A beam on two supports, having neither end supported—

We have, at once, M_i and M_k equal zero, and hence, from (B),

$$S_i = \frac{M_s}{l_i} + Q_i = 0$$
, or $M_s = -Q_i l_i \dots (312)$

And, from (C),

$$S_{s}'=0=\frac{M_{s}}{l_{s}}+Q_{s}', \text{ or } M_{s}=-Q_{s}' l_{s} \dots \dots \dots (313)$$

The shears, intermediate bending moments and deflections can now be readily obtained from the general equations.

* THE POINT OF ZERO MOMENT.

Let us take (D).

In which L is dependent upon the kind of loading in the span r. If there is no load in the span r, then

Now, if there is a point of zero moment anywhere in the span r, we can find its distance from the left support by making (314) equal zero, and solving for x_r ; doing this, we obtain

$$x_r = \frac{-M_r}{S_r} \cdot \dots \cdot (315)$$

From (B),

$$S_r = \frac{M_{r+t} - M_r}{l_r} + (Q_r \text{ in this case} = 0) \dots \dots (B)$$

Now,

$$M_m = c_m \frac{\beta_2 \beta_3 \dots \beta_{m-1}}{\beta_2'' \beta_2'' \dots \beta_{m-1}''} M_2 \dots m < r+1 \dots (70)$$

m < r+1 indicating that only those loads upon the right of the span are considered.

Substituting (B) and (70) in (315), it reduces to

$$x_r = \frac{c_r}{c_r - c_{r+1} \frac{\beta_r}{\hat{\beta}_r^n}} l_r \dots \text{ Load on the Right } \dots \dots \dots (316)$$

^{*} See "Annales des Ponts et Chaussèes," 1886, Paper No. 40, by M. Collignon.

We also have, for loads on the left,

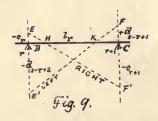
$$M_{m} = d_{s-m+2} \frac{\beta_{s-1}^{"} \beta_{s-2}^{"} \dots \beta_{m}^{"}}{\beta_{s-1} \beta_{s-2} \dots \beta_{m}} M_{s} \dots \dots (71)$$

. And, substituting (71) and (B) in (315), we obtain

$$x_{r} = \frac{d_{s-r+2}}{d_{s-r+2} - d_{s-r+\ell}} \frac{\beta_{r+\ell}''}{\beta_{r+\ell}} l_{r} . . . Load on the Left (317)$$

(316) and (317) are general equations. Knowing the points of zero moment, we can tell at once what loads must be considered to have a certain effect.

The values of x_r are very easily computed, but they can be constructed graphically as soon as c_r , c_{r+i} , etc., are known. Thus:—



Let B C represent any unloaded span. Then, B $H=x_r$ for m< r+1 or a load on the right, and B $K=x_r$ for m>r or a load on the left, if E $B=c_r$, C $F'=-c_{r+l}$ $\frac{\beta_r}{\beta_r^r}$

$$oldsymbol{F} oldsymbol{C} \! = \! oldsymbol{d}_{s-r+i} rac{eta_{r+i}^{''}}{eta_{r+i}} \; ext{and} \; oldsymbol{B} \; oldsymbol{E}' \! = \! - \; oldsymbol{d}_{s-r+2}.$$

For, from Fig. 9,

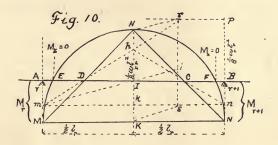
$$B H = B C \frac{E B}{E B + F' C} = \frac{c_r}{c_r + c_{r+i}} \frac{\beta_r}{\beta_r^*} l_r = x_r \dots m < r-1$$

Also,

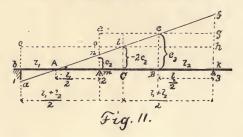
$$B \ K = B \ C \ \frac{B \ E'}{B \ E' + C \ F} = \frac{d_{s-r+2}}{d_{s-r+2} - d_{s-r+i}} \frac{\beta_{r+i}^{"}}{\beta_{r+4}^{"}} l_r = x_r \ . \ . \ m > r$$

If the span considered is loaded, then the value of x_r must be obtained from (D).

If the supports are level, and the span considered is uniformly loaded, x_r , and also the bending moments, can be found graphically.



The co-efficients c and d.



The co-efficients c and d can be found graphically, if I is considered as constant.

In Fig. 11, let b m and m k represent the first and second spans of a continuous girder. Bisect b m at A, m k at B, and b k at C. Make C l = -2 $c_2 = -2$ and draw a A l e f, a right line passing through A and l, then will B e equal, numerically, c_3 . Thus:—

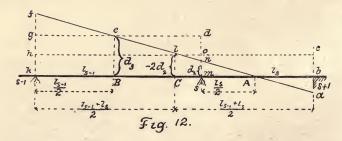
$$b \ c \ h \ k = -2 \ c_2 \ (l_1 + l_2) = -2 \ (l_1 + l_2),$$

but b c h k=f A k-a b A=f n m k,

since d e n = f e g, $f n m k = g d m k = l_2 (B e)$,

Therefore,
$$B e = -2 \frac{l_1 + l_2}{l_2} = c_3$$
.

 d_s is found in a similar manner from Fig. 12.

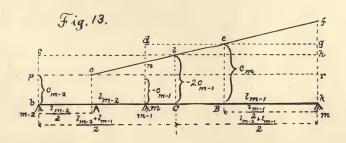


$$k B = \frac{l_{s-l}}{2}, bA = \frac{l_s}{2}, kC = \frac{l_{s-l} + l_s}{2}, C l = -2 d_2 = -2.$$

$$b\ c\ h\ k = -2\ d_{2}\ (l_{s} + l_{s-1}) = -2\ (l_{s} + l_{s-1}),$$
 $k\ g\ d\ m = (B\ e)\ l_{s-1},\ {\rm but}\ b\ c\ h\ k = k\ g\ d\ m,$
Therefore,

$$B e = -2 \frac{l_s + l_{s-1}}{l_{s-1}} = d_s.$$

Then, in general, from Fig. 13, we have, if



 $\begin{array}{l} p \; b = c_{m-2} \; \text{and} \; l \; C = -2 \; c_{m-i}, \\ c \; b \; k \; h = d \; g \; r \; d + p \; r \; k \; b = d \; g \; k \; m + p \; d \; m \; b, \\ \text{but, } g \; d \; m \; k = (B \; e) \; l_{m-i}, \; p \; d \; m \; b = c_{m-2} \; l_{m-2} \\ \text{and } c \; b \; k \; h = -2 \; c_{m-i} \; (l_{m-i} + l_{m-2}) \; . \end{array}$



Therefore,

$$B \ e = -2 \ c_{m-i} \frac{l_{m-i} + l_{m-2}}{l_{m-2}} - c_{m-2} \frac{l_{m-2}}{l_{m-i}} = c_m.$$

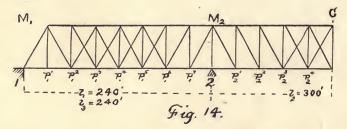
And for d_m we can write

$$B \ e = -2 \ d_{m-1} \ \frac{l_{s-m+2} + l_{s-m+3}}{l_{s-m+2}} - d_{m-2} \frac{l_{s-m+3}}{l_{s-m+2}} = d_{m}.$$

APPLICATIONS.

We will now give a complete analysis of a continuous girder, to illustrate, more particularly, the use of Table I.

Ex. 1.—Let Fig. 14 represent a continuous girder of three spans on level supports, and having a constant moment of inertia.



Let each panel be 30' in length and the loads represented as in the figure. Then s=3, r=1, 2, 3 and 4, $l_i=240'=l_3$ and $l_i=300'$.

First, look up in Table I, the values of the co-efficients $2 k-3 k^2+k^3$ and $k-k_3$, for the different values of $\frac{a}{l}=k$.

	$k_{\scriptscriptstyle f} == k_{\scriptscriptstyle S}$	$2 k_{i} - 3 k_{i}^{2} + k_{i}^{3}$	$k_i - k_i^3$	k_2	$2k_{2}$ $-3k_{2}^{2}$ $+k_{2}^{3}$ $-$	k ₂ -k ₂ ³
				0.100	0.171	-0.099
-				0.200	0 288	0.192
	0.125	0.205,078,125	0.123,046,875	0.300	0.357	0.273
	0.250	0.328,125,000	0.234,375,000	0.400	0.384	0.336
	0.375	0.380,859,375	0.322,265,625	0.500	0.375	0.375
	0.500	0.375,000,000	0.375,000,000	0.600	0.336	0 384
	0.625	0.322,265,625	0.380,859,375	0.700	0.273	0.357
	0.750	0,234,375,000	0.328,125,000	0.800	0.192	0.288
	0.875	0.123,046,875	0.205,078,125	0.900	0.099	0.171

From (A_i) , we have $M_i = 0 = M_{\lambda}$.

$$M_{z} = \frac{-c_{z}}{d_{x} l_{i}} \left\{ \frac{A_{z} d_{z} + B_{i} d_{z}}{A_{z} d_{z} + B_{z} d_{i}} \right\} + \frac{-d_{z}}{c_{x} l_{z}} \left\{ A_{i} c_{i} + B_{i} c_{z} \right\}$$

$$M_s = rac{-c_s}{d_s \ l_s} \left\{ \ A_s \ d_z + B_s \ d_s \
ight\} \ + rac{-d_z}{c_s \ l_s} \left\{ egin{matrix} A_s \ c_i + B_s \ c_z \ A_z \ c_z + B_z \ c_s \ \end{pmatrix}
ight.$$

From (p_i) and (r_i) , we find that, since $c_i=0$ and $c_2=1$, $c_3=-3.6$ and $c_4=+14.95$.

Also,

$$d_1 = 0$$
 and $d_2 = 1$, $d_3 = -3.6$ and $d_4 = +14.95$.

Finding the values of A_i , A_2 , B_i , B_2 , etc., from (i_i) and (j_i) , and substituting them, with the values of c and d above, in the moment equations, we obtain,

TABLE a-Values of M_2 .

```
 \begin{cases} -7.111,198 \ P_1^1 \\ -13.545,140 \ P_1^2 \\ -18.624,568 \ P_3^3 \\ -21.672,225 \ P_1^4 \\ -23.010,833 \ P_5^5 \\ -18.963,196 \ P_1^6 \\ -11.851,998 \ P_7^7 \end{cases} 
 - P_1^1 (0.123,046,875)
 -P^{\frac{5}{2}}(0.234.375.000)
 -P_1^3 (0.322,265,625)
 -P_1^4(0.375,000,000)
 -P_1^5 (0.380,859,375)
-P_1^6 (0.328,125,000)
    - P<sup>7</sup> (0.205,078,125)
                                                                                            + P_{2}^{1}(0.099)
                                                                                                                                                 ( -- 12.958,194 Pt
 -P_{2}^{1}(0.171)
 \begin{array}{l} -P_{\frac{1}{2}}^{1}\left(0.171\right) \\ -P_{\frac{3}{2}}^{2}\left(0.288\right) \\ -P_{\frac{3}{2}}^{3}\left(0.357\right) \\ -P_{\frac{1}{2}}^{4}\left(0.384\right) \\ -P_{\frac{5}{2}}^{2}\left(0.375\right) \\ -P_{\frac{5}{2}}^{6}\left(0.336\right) \\ -P_{\frac{7}{2}}^{7}\left(0.273\right) \\ -P_{\frac{5}{2}}^{8}\left(0.192\right) \\ -P_{\frac{5}{2}}^{8}\left(0.192\right) \\ -P_{\frac{5}{2}}^{8}\left(0.1992\right) \\ -P_{\frac{5}{2}}^{8}\left(0.1992\right) \\ -P_{\frac{5}{2}}^{8}\left(0.1992\right) \end{array} \right\} 
                                                                                                                                    \begin{array}{c} -12.935,194 \text{ Fz} \\ -21.190,636 \text{ Pz} \\ -25.389,634 \text{ Ps} \\ -26.247,494 \text{ Ps} \\ -24.456,525 \text{ Ps} \\ -20.709,033 \text{ Ps} \\ -15.697,327 \text{ Ps} \\ -10.113,715 \text{ Ps} \\ -4.650 \text{ 503 Ps} \end{array}
                                                                                                                                    -4.650,503 P_2^9
 + P_{2}^{1} (0.205,078,125)
                                                                                                                                                       + 3.292.221 Pt
                                                                                                                                                      + 5.267,554 P_3^2
 + P_{3}^{2}(0.328,125,000)
                                                                                                                                                      + 6.114,126 P3
 + P_{3}^{3} (0.380.859.375)
                                                                                                                                                   +6.020,061 P<sup>3</sup>
                                                              16.0535 = \dots \dots
 + P^{\frac{1}{2}} (0.375.000,000) 
                (0.322,265,625)
                                                                                                                                                      + 5.173,491 P_3^5
 + P_3^6 (0.234,375,000)
                                                                                                                                                      + 3.762,365 P
 + P_{5}^{9} (0.123,046,875)
                                                                                                                                                      + 1.975,333 P_3^7
```

TABLE a-Continued.

VALUES OF M3.

```
+ P_1^1 (0.123,046,875)
                                                                                                                                                                                                                                                                                         + 1.975,333 P
+ P_1^{\frac{1}{2}} (0.234, 375,000)
                                                                                                                                                                                                                                                                                          + 3.762,365 P_1^2
+ P_1^{\frac{1}{3}} (0.322, 265, 625)
                                                                                                                                                                                                                                                                                       + 5.173,491 P_1^3
+ P_1^4 (0.375,000,000)
                                                                                                                      16.0535 = ...
                                                                                                                                                                                                                                                                                      + 6.020,061 P_1^4
+ P_1^{\frac{1}{5}} (0.380, 859, 375)
                                                                                                                                                                                                                                                                                         + 6.114,126 P_1^5
+ P_1^{\frac{1}{6}} (0.328, 125,000)
                                                                                                                                                                                                                                                                                         + 5.267,554 P_1^6
                                                                                                                                                                                                                                                                                          + 3.292,221 P7
 + P_7^7 (0.205,078,125)
 + P_2^1 (0.171)
                                                                                                                                                                                  -P_2^1(0.099)
                                                                                                                                                                                                                                                                                   - 4.650,503 P<sub>2</sub>
                                                                                                                                                                                                                                                                                   -10.113,715 P_2^2
  + P_{\frac{2}{3}}(0.288)
                                                                                                                                                                                            P2 (0.192)
                                                                           \begin{array}{c} -\frac{1}{2} & (0.192) \\ -\frac{1}{2} & (0.273) \\ -\frac{1}{2} & (0.336) \\ -\frac{1}{2} & (0.336) \\ -\frac{1}{2} & (0.336) \\ -\frac{1}{2} & (0.375) \\ -\frac{1}{2} & (0.384) \\ -\frac{1}{2} & (0.288) \\ -\frac{1}{2} & (0.288) \\ -\frac{1}{2} & (0.287) \\ -\frac
+ P_{2}^{3} (0.357) 
+ P_{2}^{4} (0.384)
                                                                                                                                                                                                                                                                             = \{ -24.456,525 \text{ P}_{2}^{5} \}
+ P_{2}^{5} (0.375)  + P_{2}^{6} (0.336) 
                                                                                                                                                                                                                                                                              - 26.247,494 P
+ P_{2}^{7} (0.330) + P_{2}^{7} (0.273)
                                                                                                                                                                                                                                                                                   - 25.389,634 P
                                                                                                                                                                                                                                                                                 -21.190,636 P_2^8
 + P_2^8 (0.192)
+ P_2^{9} (0.099)
                                                                                                                                                                                  -P_{2}^{5}(0.171)
                                                                                                                                                                                                                                                                               -15.958,194 P_2^9
    - P<sup>1</sup><sub>2</sub> (0.205,078,125)
                                                                                                                                                                                                                                                                                             - 11.851,998 P<sub>3</sub>
  -P_3^2(0.328,125,000)
                                                                                                                                                                                                                                                                                             - 18.963,196 P<sup>2</sup><sub>3</sub>
\begin{array}{l} - P_3^3 \ (0.380,859,375) \\ - P_3^4 \ (0.375,000,000) \end{array}
                                                                                                                                                                                                                                                                                  - 22.010,853 Ps
                                                                                                                                                                                                                                                                                  \begin{array}{l} -21.672,225 \text{ P}_{3}^{4} \\ -18.624,568 \text{ P}_{5}^{5} \\ -13.545,140 \text{ P}_{6}^{6} \end{array}
\begin{array}{c} -P_{3}^{5} (0.375,000,000) \\ -P_{3}^{5} (0.322,265,625) \\ -P_{3}^{6} (0.234,375,000) \end{array}
-P_{3}^{7}(0.123,046,875)
                                                                                                                                                                                                                                                                                      — 7.111,198 P<sub>3</sub>
```

Compare the values of M_2 and M_3 , and notice that the co-efficients of P_3^7 , P_3^6 , etc., of M_3 , are the same as the co efficients of P_1^r , P_2^r , etc., of M_2 , as they should be, owing to the symmetry of the girder.

TABLE b.

The next step is the deduction of S from (B) and (C).

		FIRST	SPAN.	SECONI	SPAN.	THIRD SPAN.		
	-	S_i	$S_2^{'}$	S_2	S_3'	S_3	S_4	
FIRST SPAN.	$\begin{array}{c} P_1^1 \\ P_2^2 \\ P_3^3 \\ P_1^4 \\ P_5^5 \\ P_1^6 \\ P_7^7 \end{array}$	$\begin{array}{c} +0.845,370 \\ +0.693,562 \\ +0.547,397 \\ +0.409,699 \\ +0.283,288 \\ +0.170,986 \\ +0.075,616 \end{array}$	+0.154,630 +0.306,438 +0.452,603 +0.590,301 +0.716,712 +0.829,014 +0.924,384	+0.030,288 +0.057,691 +0.079,327 +0.092,307 +0.093,749 +0.080,769 +0,050,480	Same as S ₂ with—sign	-0.008,230 -0.015,676 -0.021,556 -0.025,083 -0.025,475 -0.021,948 -0.013,717	Same as \$3 with+sign	
SECOND SPAN.	$\begin{array}{c} P_{2}^{1} \\ P_{2}^{2} \\ P_{3}^{2} \\ P_{4}^{2} \\ P_{5}^{2} \\ P_{5}^{2} \\ P_{7}^{2} \\ P_{9}^{2} \\ \end{array}$	$\begin{array}{c} -0.053,992 \\ -0.088,294 \\ -0.105,790 \\ -0.101,902 \\ -0.101,902 \\ -0.086,287 \\ -0.065,405 \\ -0.042,140 \\ -0.019,377 \end{array}$	Same as \$1 with+sign	+0.927,692 +0.836,923 +0.732,307 +0.618,461 +0.590,000 +0.381,538 +0.267,692 +0.163,076 +0.072,307	+0.072,307 +0.163,076 +0.267,692 +0.381,538 +0.500,000 +0.618,461 +0.732,307 +0.836,923 +0.927,692	+0.019,377 +0.042,140 +0.065,405 +0.086,287 +0.101,902 +0.109,364 +0.105,790 +0.088,294 +0.053,992	Same as \$3 with—sign	
THIRD SPAN.	$\begin{array}{c} P_{3}^{1} \\ P_{3}^{2} \\ P_{3}^{3} \\ P_{3}^{4} \\ P_{5}^{5} \\ P_{7}^{6} \end{array}$	$\begin{array}{c} +0.013,717 \\ +0.021,948 \\ +0.025,475 \\ +0.025,083 \\ +0.021,556 \\ +0.015,676 \\ +0.008,230 \end{array}$	Same as \$1 with-sign	$\begin{array}{c} -0.050,480 \\ -0.086,769 \\ -0.093,749 \\ -0.092,307 \\ -0.079,327 \\ -0.057,691 \\ -0.030,288 \end{array}$	Same as S ₂ with+sign	$\begin{array}{c} +0.924,384 \\ +0.829,014 \\ +0.716,712 \\ +0.590,301 \\ +0.452,603 \\ +0.306,438 \\ +0.154,630 \end{array}$	$\begin{array}{c} +0.075,616 \\ +0.170,986 \\ +0.283,288 \\ +0.409,699 \\ +0.547,397 \\ +0.693,562 \\ +0.845,370 \end{array}$	

The computation of the above table is very simple and easy. Thus:—

 $S_i = \frac{M_2}{l_i} + P_i (1-k_i)$, and numerically becomes, for loads in the first span,

$$S_{i} = \left(\frac{-7.111 + 210}{240} + \frac{210}{240}\right) P_{i}' = +0.845 P_{i}'.$$

$$S_{i} = \left(\frac{-13.545}{240} + \frac{180}{240}\right) P_{i}' = +0.693 P_{i}', \text{ etc., etc.}$$

For loads in the second and third spans, we have merely M_2 , or the moment over the second support divided by the length of the first span.

$$S_z = rac{M_s - M_z}{l_z} + P_z$$
 (1- k_z), for loads in the second span, and $S_z = rac{M_s - M_z}{l_z}$, for loads in the other spans. $S_s = rac{-M_z}{l_s} + P_s$ (1- k_s), for loads in the third span, and

 $S_3 = \frac{-M_2}{l_3}$, for loads in the other spans.

It is necessary to compute only S_i , S_2 and S_3 to fill out the table of shears, since $S_i + S_2' = 0$ or P_i , $S_2 + S_3' = 0$ or P_2 , etc., but it is better to compute S_i , S_2' and S_3' , as a check.

TABLE c-VALUES OF Mz.

E	T	D	C	T	2	D 3	IN.

0	1	2	3	4	õ	в	7_	8	9
		M_x^{i}	$M_x^{'2}$	$M_{x_i}^{i_3}$	$M_x^{f_A}$	$M_x^{'5}$	$M_x^{'e}$	$M_x^{'7}$	
1 2 3 4 5 6 7	$\begin{array}{c} P_1^1 \\ P_2^2 \\ P_3^3 \\ P_4^4 \\ P_5^5 \\ P_6^6 \\ P_7^7 \end{array}$	+ 25.361 + 20.806 + 16.421 + 12.290 + 8.498 + 5.129 + 2.268	+ 20.722 + 41.613 + 32.843 + 24.581 + 16 997 + 10.259 + 4,536	+ 16.083 + 32.420 + 49.265 + 36.872 + 25.495 + 15.388 + 6.805	+ 11.444 + 23.237 + 35.687 + 49.163 + 33.994 + 20.518 + 9.073	+ 6.805 + 14.044 + 22.109 + 31.554 + 42.493 + 25,647 + 11.342	+ 2.166 + 4.851 + 8.531 + 13.745 + 20.991 + 30.777 + 13.610	$\begin{array}{r} -2.472 \\ -4.341 \\ -5.046 \\ -3.963 \\ -0.509 \\ +5.907 \\ +15.879 \end{array}$	P ₃ ⁷ P ₆ ⁸ P ₃ ⁸ P ₃ ⁴ P ₃ ⁸
8 9 10 11 12 13 14 15 16	$\begin{array}{c} P_{2}^{1} \\ P_{2}^{2} \\ P_{2}^{3} \\ P_{2}^{4} \\ P_{2}^{5} \\ P_{2}^{6} \\ P_{2}^{7} \\ P_{2}^{6} \\ P_{2}^{9} \end{array}$	- 1.619 - 2.648 - 3.173 - 3.280 - 3.057 - 2.588 - 1.962 - 1.264 - 0.581	- 3.239 - 5.297 - 6.347 - 6.561 - 6.114 - 5.177 - 3.924 - 2.528 - 1.162	- 4.859 - 7.916 - 9.521 - 9.842 - 9.171 - 7.765 - 5.886 - 3.792 - 1.743	- 6.479 - 10.595 - 12.694 - 13.123 - 12 228 - 10.354 - 7.848 - 5.056 - 2.325	- 8,698 - 13,244 - 15,868 - 16,404 - 15,285 - 12,943 - 9,810 - 9,321 - 2,906	- 9.718 -15 892 -19.042 -19.685 -18.342 -15.531 -11.772 -7.585 -3.487	- 11.338 - 18.541 - 22.215 - 22.966 - 21.399 - 18.120 - 13.735 - 8.849 - 4.0 9	P2 P2 P2 P2 P2 P2 P2 P2 P2 P2 P2 P2 P2 P
17 18 19 20 21 22 23	$\begin{array}{c} P_{3}^{1} \\ P_{3}^{2} \\ P_{3}^{3} \\ P_{3}^{4} \\ P_{5}^{5} \\ P_{6}^{5} \\ P_{7}^{7} \end{array}$	+ 0.411 + 0.658 + 0.764 + 0.752 + 0.646 + 0.470 + 0.246	+ 0.823 + 1.316 + 1.528 + 1.504 + 1.293 + 0.940 + 0.493	$\begin{array}{c} +\ 1.234 \\ +\ 1.975 \\ +\ 2.292 \\ +\ 2.257 \\ +\ 1.940 \\ +\ 1.410 \\ +\ 0.740 \end{array}$	+ 1.646 + 2.633 + 3.057 + 3.009 + 2.586 + 1.881 + 0.987	+ 2.057 + 3.292 + 3.821 + 3.762 + 3.233 + 2.351 + 1.234	+ 2.469 + 3.950 + 4.585 + 4.514 + 3.880 + 2.821 + 1.481	+ 2.880 + 4.609 + 5.349 + 5.267 + 4.527 + 3.291 + 1.728	P ₁ P ₆ P ₅ P ₁ P ₁ P ₁ P ₁ P ₁
		$_{_{_{_{_{_{_{_{_{_{_{_{_{_{_{_{_{_{_$	$\stackrel{_3}{M_x^6}$	$\stackrel{_3}{M_x^5}$	$\stackrel{_3}{M_x^4}$	$\stackrel{_{_{_{_{_{_{_{_{_{_{_{_{_{_{_{_{_{}}}}}}$	$\stackrel{_{_3}}{M_x^2}$	M_x^{i}	

THIRD SPAN.

TABLE c-Continued.

SECO	ND	SPA	N.

199		The second second	-							
1	2	3	4	5	6	7	8	9	10	11
-	M_x^2	M_x^2	M_x^2	$\mathring{M_x^2}$	M_x^{25}	M_x^{s}	$\hat{M_x^2}$	M_x^2	M_x^{2g}	
$2 P_{1}^{2}$	- 6.202 - 11.824 - 16.244	$-10\ 103$	- 8,382	-6.662	-4.941	-3.220		+ 0.220	+1.941	P_3^6
4 P ⁴ 5 P ⁵	-18.903 -19.198 -16.540	- 16.133 - 16.385	- 13.364 - 13.573	-10.595 -10.760	-7826 -7.948	- 5.056 - 5.136	- 2.287 - 2.323	$^{+}$ 0.481 $^{+}$ 0.498	$^{+}$ 3.250 $^{+}$ 3.301	$ \begin{array}{c} P_{3}^{4} \\ P_{3}^{3} \\ P_{3}^{2} \end{array} $
$7P_1^7$	- 10.337	- 8.823	- 7.308 	- 5.794	- 4.279	- 2.765	- 1.251	+ 0.263	+ 1.777	P_3^{i}
$\frac{9 \text{ P}_2^2}{10 \text{ P}_2^3}$	+ 14.872 + 3.917 - 3.420	+29.024 +18.548	+24.132 +40.517	$+\ 19.240 \\ +\ 32\ 487$	$+\ 14.347 \\ +\ 24.456$	$+\ 9.355 +\ 16.425$	$^{+}$ 4.563 $^{+}$ 8.394	$-0.329 \\ +0.364$	-5.221 -7.666	$\frac{P_2^8}{P_2^7}$
	- 7.693 - 9,456	+ 5.543	+ 20.543	+ 35.543	+ 50.543	+ 35.543	+ 20 543	+ 5.543	- 9.456	
	M_x^{2g}	M_x^2	$M_x^{2\gamma}$	M_x^{26}	M_x^{25}	M_x^2	$M_x^{\frac{2}{3}}$	$M_x^{\frac{2}{2}}$	M_x^2	

SECOND SPAN.

 $\mathbf{M}_{x}^{i} = \mathbf{S}_{i} \ \mathbf{x}_{i} - \mathbf{P}_{i} \ (\mathbf{x}_{i} - \mathbf{a}_{i})$ for loads in the first span, and $\mathbf{M}_{x}^{i} = \mathbf{S}_{i} \ \mathbf{x}_{i}$ for loads in the other spans. $\mathbf{M}_{x}^{3} = \mathbf{S}_{3} \ \mathbf{x}_{3} - \mathbf{P}_{3}$ ($\mathbf{x}_{3} - \mathbf{a}_{3}$) for loads in the third span, and $\mathbf{M}_{x}^{3} = \mathbf{S}_{3} \ \mathbf{x}_{3}$ for loads in the other spans.

If there were no equal spans, then Table c would have the number of apices squared computations, making its formation tedious, if many spans were considered, but the great worth of the table, when once computed, overbalances the hard labor in its formation.

Having Table c before us, we can tell at once which of the apices must be loaded, to produce a certain result, in any particular chord member.

MAXIMUM MOMENTS.

Table (c) enables one to find the maximum bending moments for any chord piece, with comparatively little labor.

Dead load-

For the dead, or static loads, of the structure, the maximum bending moment for any chord member is found by taking the algebraic sum of the quantities in the proper column of Table (c) (each co-efficient is, of course, multiplied by its P).

For example, suppose the maximum bending moment at the second panel of the first span, or, better, the second panel point of the first span, is desired: Multiply each co-efficient in column 3 of Table (c), by its proper P, and take the algebraic sum of the products.

Live load-

For the live or moving load, the maximum negative moment at any panel point is found by taking the sum of the negative co-efficients (multiplied by their proper P's) in the proper column of Table (c), and vice versa for the maximum positive bending moment.

For example, suppose the maximum negative bending moment at the second panel point is desired: Take the sum of the negative products in column 3 of Table (c). For positive moment, take the sum of the positive products.

If the P's are equal, the work is somewhat easier, as the co-efficients can be at once summed, and then multiplied by the common value of P.

The maximum bending moments over the supports are obtained in the same manner from Table (u).

MAXIMUM SHEAR.

The maximum shear is obtained from Table (b).

.Dead load-

The maximum shear at any panel point is found by:

First.—Taking the algebraic sum of the co-efficients (multiplied by their P's) of all the loads outside of the span in which the apex considered lies.

Second.—Considering the span in which the apex lies alone, and using for the left reactions the co-efficients as found in the column giving the values of S.

For example, take the second apex of the first span. The maximum shear is found: First, by taking the algebraic sum of the co-efficients in column S_i for the second and third spans. Second, by taking the sum of the co-efficients opposite P_i^s , P_i^s to P_i^r , inclusive, and decreasing it by the co-efficient opposite P_i^t in column S_i^s .

Live load-

For positive or negative shear:

First.—Take the sum of the positive or negative co-efficients (multiplied by their P^{s}) for spans in which the apex does not lie.

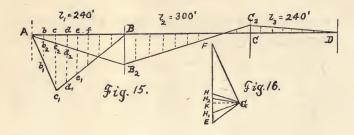
Second.—Consider the span in which the apex lies alone, and use for left reactions the positive or negative co-efficients (multiplied by their P's) as given in the proper column of Table (b).

For example, find the maximum positive shear at the second apex of the first span:

First.—Take sum of co-efficients (multiplied by their P's) in column S_i , opposite third span.

Second.—Take sum of co efficients (multiplied by their P^{s}), opposite P_{i}^{s} to P_{i}^{s} , inclusive.

Knowing the maximum shears and bending moments, the maximum stresses are then easily obtained.



The following is a graphical solution of the same example. Tables (b) and (c) can be filled by means of graphics, as can also Table (a), by using Prof. Greene's Method of Area Moments, which is fully explained in Part II of Greene's Trusses and Arches.

Let it be supposed that Table (a) has been filled, either by computation or by Greene's Method of Area Moments; we have, then, the bending moments over the supports, or the "pier ordinates." Consider only a single concentration at the second apex of the first span, and call it P^2 ; then, if the first span were discontinuous, the moment polygon $A c_i B$ would enable us to determine the bending moments at other apices of the span; but the girder is not discontinuous, and the load P^2 induces a negative moment over the second support, which may be designated by $B B_2$, then the other negative moments are given by the ordinates in the triangle or polygon A B B. The difference of the ordinates of the two polygons determines the magnitude and kind of bending moment at each apex of the first span. The load P_i^2 induces a positive moment over the second support, which may be designated by C C_2 , then the ordinates between the lines B C and B_2 C_2 will give us the moments at the apices of the second span. The

moments in the third span are given by the ordinates between the lines C_2 D and C D.

Having given an outline of the method, we will now proceed to give it in detail.

First.—Assume some value for P_i^2 , as unity, ten or one hundred.

Second.—Form the stress diagram, Fig. 16, assuming for the pole distance some value as unity, ten or one hundred. Assume the pole in such a position, that the closing line \boldsymbol{A} \boldsymbol{B} to the equilibrium polygon \boldsymbol{A} \boldsymbol{c}_i \boldsymbol{B} shall be horizontal, and construct the equilibrium polygon \boldsymbol{A} \boldsymbol{c}_i \boldsymbol{B} .

Third.—From Table (a), we find the bending moment over the second support to be — $13.54 P_i^2$; now, if $P_i^2 = 1$, and the pole distance = 1, then the ordinate $B B_2 = 13.54$, laid off to the scale of A B. Draw $A B_2$. Scale the ordinates $b_i b_2$, $c_i c_2$, etc., and place the results in Table (c), column 3, opposite P_i^i , P_2^2 , etc. The results will be found to agree with those computed.

Fourth.—From Table (a), we find that $C C_2 = +3.762$ P_i^2 ; hence, lay off $C C_2 = 3.76$ to scale of A B, and draw $B_2 C_2$ and $C_2 D$ and fill out the remainder of column 3, Table (c), by scaling the ordinates between these lines and the horizontal.

In like manner, each column of Table (c) can be filled in a comparatively short time.

To fill Table (b), proceed as follows:

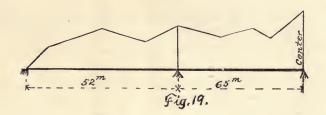
First. — Draw G H, Fig. 16, parallel to A B_2 , then E H= S_1 and F H= S_2 .

Second.—Draw G H_i parallel to B_2 C_2 , then K $H_i = S_2$ $= -S_3$.

Third.—Draw G H_z parallel to C_z D, then K H_z = $-S_x$ =+ S_x '.

Proceed in like manner with each load.

* EXAMPLE 2—VARIABLE I.



$$l_s = 52$$
, $l_s = 65$, $l_s = 65$, $l_s = 52$. $s = 4$. $r = 1$, 2, 3 and 4. $m = 1$, 2, 3, 4 and 5.

Determine the value of M_2 .

From (A),

$$M_{z} = \frac{-1}{(d_{s} \beta_{s}^{"})} \left\{ \begin{matrix} (A_{s} + X_{s}^{'} + X_{s}^{"}) \ d_{s} \\ (A_{s} + X_{s}^{'} + X_{s}^{"} + B_{s}) \ d_{s} \\ (A_{s} + X_{s}^{'} + X_{s}^{"} + B_{s}) \end{matrix} \right\} + \frac{-d_{s} \beta_{s}^{"} \beta_{s}^{"}}{(c_{s} \beta_{s}) \beta_{s} \beta_{s}} B_{s}$$

In which the several terms have the following values:

^{*} See page 131, Allyemeine Theorie und Berechnung der continuirlichen und einfachen Träger, by J. Jacob Weyrauch, Ph. D.

First Span.

	$l_i = 52$	$e_l = e_s$	$I_l = I_s$.	
ℓ_v			\triangle_v	$\triangle_v \frac{e_v^3}{l_I}$
$e_{1} = e_{2} = e_{2} = e_{3} = e_{4} = e_{5} = e_{6} = e_{7} = e_{8} = e_{8} = e_{7} = e_{8} = e_{7} = e_{8} = e_{7} = e_{8} = e_{8} = e_{7} = e_{8} = e_{8} = e_{7} = e_{7} = e_{8} = e_{7} = e_{7} = e_{8} = e_{7} = e_{7$	10.50 13.75 31.25 34.50 43.50 46.00 48.00 50.00		$\begin{array}{c} +\ 0.300 \\ +\ 0.337 \\ -\ 0.337 \\ -\ 0.300 \\ +\ 0.433 \\ +\ 0.289 \\ +\ 0.201 \\ +\ 0.243 \end{array}$	$ \begin{array}{r} + & 6 \\ + & 17 \\ - & 198 \\ - & 237 \\ + & 685 \\ + & 541 \\ + & 427 \\ + & 584 \\ \hline + & 1825 \\ \end{array} $
		1895		

$F_i'' = + \frac{1825}{52} = + 35.0961$

Second Span.

$l_2 = 65.$ $e_l = e_{l3}.$ $I_l = I_{l3}.$							
ℓ_v	\triangle_v	△v ·	$\frac{l^3 - (l - e_v)^3}{l}$	$\left \triangle_v \; \epsilon_i \right $	$\left(2 - \frac{e_v}{l}\right)$	$\triangle v$	$\frac{e_v^3}{l}$
$e_i = 1.50$	$\triangle_i = -0.256$		73		2		0
$e_z = 3.00$	$\triangle_2 = -0.222$		124		6		0
$e_s = 4.75$	$\triangle_s = -0.311$	_	267	_	20		1
$e_{4} = 6.75$	$\triangle_{s} = -0.467$		553		59		2
$e_{s} = 19.25$	$\triangle_{5} = +0.467$	+	1285	+	417	+	51
$e_{\theta} = 23.50$	$\triangle_{\theta} = +0.311$	+	-972	+	391	+	62
$e_7 = 42.75$	$\triangle_7 = -0.311$		1261		957	_	374
$e_s = 47.00$	$\triangle_s = -0.467$		1931		1603		746
$e_{\mathfrak{g}} = 5625$	$\triangle_{\mathfrak{g}} = +0467$	+	1968	+	1876	+ :	1278
$e_{10} = 58.75$	$\triangle_{10} = +0.311$	+	1313	+	1279	+	970
$e_{II} = 60.75$	$\triangle_{11} = +0.222$	+	938	+	927	+	765
$e_{12} = 62.25$	$\triangle_{12} = +0.167$	+	706	+	701	+	620
$e_{i3} = 63.50$	$\triangle_{13} = +0.166$	+	701	+	700	+	657
/			+3674	+	3641		3280
$\frac{+3674}{65} = +56.524 = F_2. \qquad \frac{+3644}{65} = +56.061 = F_2'.$							
$\frac{+3280}{65} = +50461 = F_z''.$							

Third Span.

	$l_3 = 65.$	$e_l = \epsilon$	e_{is} . $I_i =$	$=I_{i3}$.			
\mathscr{C}_v	△v	$\left \triangle_v \right ^2$	$\frac{l-(l-e_v)^3}{l}$	$\triangle_v e$	$-2 \frac{e_v}{l}$	$\triangle v$	$\frac{\mathscr{C}_{v}^{3}}{l}$
$e_i = 1.50$ $e_2 = 2.75$ $e_3 = 4.25$	$\begin{vmatrix} \triangle_{i} = -0.154 \\ \triangle_{2} = -0.155 \\ \triangle_{3} = -0.206 \end{vmatrix}$		44 80 160	_	1 3 - 11		0 0
$e_4 = 6.25$ $e_5 = 8.75$ $e_6 = 18.00$ $e_7 = 22.25$	$ \begin{array}{l} $	_ + +	319 644 1138 874	+++	32 90 343 331	++	1 4 39 49
$e_s = 41.50$ $e_g = 45.75$ $e_{to} = 58.25$ $e_{tt} = 60.25$	$ \begin{array}{l} $	+++	1163 1782 1827 1220	+++	858 1443 1775 1202	- + 1 +	318 638 316 972
$e_{i2} = 62.00$ $e_{i3} = 63.50$		++	- 870 1005	+	865 1004	++	755 937
$\frac{-}{+2742} =$	$+42.215 = F_s$.	1 +	$ \begin{array}{r} 2742 \\ +3082 \\ \hline 65 \end{array} $		+47.4		F_{3}^{\prime} .
	$\frac{+3107}{65}$ =	= +	47.800 =	$F_{s}^{"}$.			

Fourth Span.

		$l_{\star}=52.$	$e_l =$	$=e_8.$	$I_t = I_s$		
e_{ϵ}	,	_	V	$\triangle_v \frac{l^3-}{}$	$\frac{-(l-e_v)^3}{l}$	$\triangle_v e_v^2 (3$	$-2\frac{e_v}{l}$
$e_{1} = e_{2} = e_{3} = e_{4} = e_{5} = e_{6} = e_{7} = e_{8} = e_{8} = e_{8$	2.00 4.00 6.00 8.50 17.50 20.75 38.25 41.50		0.110 0.095 0 134 0.200 0.143 0.151 0.151 0.143	 + + 	33 55 111 224 274 320 401 384	1 + +	1 4 13 39 102 143 337 346
				-	— 614		495
$\frac{-614}{52}$	= -1	1.807=F	۲ 4		$\frac{495}{52} =$	- 9,51	$9=F_{\lambda}'$

Collecting the values of
$$F$$
, we have

$$F_{s} = + 56.523 \quad F_{s}' = + 56.061 \quad F_{s}'' = + 50.461 \quad F_{s} = + 42.215 \quad F_{s}' = + 47.415 \quad F_{s}'' = + 47.800 \quad F_{s}'' = + 47.8$$

can write

$$\Sigma P_{i} l_{i} (1-k_{i}) = \frac{1}{2} w_{i} l_{i}^{2} = \frac{1}{2} \cdot 6.72704 = \dots 9,058.4$$

$$\Sigma P_2 l_2 (1-k_2) = \frac{1}{2} w_2 l_2^2 = \frac{1}{2} 6.74225 = \dots 14,153.7$$

$$\Sigma P_s l_s (1-k_s) = \frac{1}{2} \hat{w}_s l_s^2 = \frac{1}{2} \cdot 2.24225 = \dots 4,647.5$$

$$\Sigma P_{s} l_{s} (1-k_{s}) = \frac{1}{2} w_{s} l_{s}^{2} = \frac{1}{2} 6.72704 = \dots 9,058.4$$

Also,

$$\Sigma P_r (e_v - a_r)^s = \int_0^{e_v} w_r d a_r (e_v - a_r)^s = \frac{1}{4} w_r e_v^4$$

$$e_v$$

$$\sum_{i=0}^{n} P_r (e_v - a_r)^2 = \int_{0}^{e_v} w_r d a_r (e_v - a_r)^2 = \frac{1}{3} w_r e_v^3$$

Substituting these values in (f) and (g), they become

$$H_{r} = \sum_{v=l_{r}}^{r} \sum_{v=l_{r}}^{r} \left\{ \frac{w_{r} e_{v}^{A}}{4 l_{r}} + \frac{3 (l_{r} - e_{v})}{l_{r}} \frac{w_{r} e_{v}^{3}}{3} \right\}$$

$$v = 1$$

$$= \sum_{v=l_{r}}^{r} \sum_{v=l_{r}}^{r} \frac{w_{r} e_{v}^{3}}{4} \left\{ 4 - 3 \frac{e_{v}}{l_{r}} \right\}$$

$$v=1$$
 $H_r'=\sum_{v=l_r}^r \left\{ \frac{w_r e_v^4}{4 \ l_r} - \frac{3}{l_r} \frac{e_v}{3} \right\}$
 $v=l_r$
 $v=1$
 $v=l_r$
 $v=l_r$
 $v=l_r$

VALUES OF \boldsymbol{H} AND \boldsymbol{H}' .

	l_{z}	l_2		l_s		· · · · · · · · · · · · · · · · · · ·
$ \begin{array}{c} v \\ \hline 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \\ 9 \\ 10 \\ 11 \\ 12 \\ 13 \\ \end{array} $	$ \begin{array}{c c} & & \\$	$\begin{array}{c} + & 10365 \\ + & 11767 \\ - & 49249 \\ - & 88767 \\ + & 117518 \\ + & 81258 \\ + & 59538 \\ + & 45396 \end{array}$	+1459 -15980 -35059 $+71927$ $+57000$ $+46518$	$\begin{array}{l} + & 8003 \\ + & 9464 \\ - & 43059 \\ - & 78303 \\ + & 112239 \\ + & 77063 \\ + & 55895 \end{array}$	$ \begin{array}{r} + 699 \\ + 1090 \\ - 13188 \\ - 29183 \\ + 76693 \\ + 58589 \end{array} $	$ \begin{array}{c c} \widehat{s} & \widehat{p} \\ \widehat{s} & \widehat{p} \\ - & 23 \\ - & 106 \\ - & 431 \\ + & 2292 \\ + & 3781 \\ - & 15153 \\ - & 16412 \end{array} $
	$H_{i}^{'}$	H_{z}	$H_{z}^{'}$	$H_{\scriptscriptstyle 3}$	$H_{\mathfrak{A}}^{'}$	H_{ι}

 $-67777 w_4 + 58147 w_2 - 155203 w_2 + 51270 w_3 - 150762 w_3 - 6514 w_4$

$$H_s = +389,584.9.$$
 $H_t' = -454,105.9.$ $H_s = +112,794.0.$ $H_s' = -1,039,860.1.$ $H_s' = -331,676.4.$

-311,716,223 E'

$$X_s = \begin{cases} +F_s' & \Sigma P_s \ l_s \ (1-k_s) \ \theta_s = \\ -47.41 \ (4647.5) \ 560 \ E' = -123,389,266 \ E' \\ -2 & F_s'' & \Sigma P_s \ l_s \ (1-k_s) \ \theta_s = \\ -2 \ (50.46) \ (14153.7) \ 520 \ E' = -742,763,530 \ E' \\ +H_s \ \theta_s = +112,794.0 \ (560) \ E = +63,164,640 \ E' \\ -H_s' \ \theta_s = +1,039,860.1 \ (520)E' = +\frac{540,727,252}{262,260,904} \ E' \\ -262,260,904 \ E' \\ -262,260,904 \ E' \\ -27,80 \ (4647.5) \ 240 \ E' = -106,632,240 \ E' \\ -27,80 \ (4647.5) \ 240 \ E' = -106,632,240 \ E' \\ -27,80 \ (4647.5) \ 240 \ E' = -22,694,776 \ E' \\ -27,80 \ (4647.5) \ 240 \ E' = +\frac{79,602,336}{4,881,977} \ E' \\ -27,80 \ (4647.5) \ 240 \ E' = -22,694,776 \ E' \\ -27,80 \ (47,80) \ (47,80) \ (47,80) \ (47,80) \ (47,8$$

From
$$(p)$$
, since $c_i=0$, and $c_2=1$,

$$c_{3} = \frac{-2 \beta'_{2}}{\beta_{2}} = -3.557.$$
 $c_{4} = \frac{-2 c_{3} \beta'_{3} - \beta''_{2}}{\beta_{3}} = +29.33.$ $c_{5} = \frac{-2 c_{4} \beta'_{4} - c_{3} \beta''_{3}}{\beta_{4}} = \frac{-2.590,360 \text{ E}'}{\beta_{4}}$

From (r), since $d_1=0$, and $d_2=1$,

$$d_{s} = \frac{-\beta'_{s}}{\beta''_{s}} = -1.524. \qquad d_{s} = \frac{-2 d_{s} \beta'_{s} - \beta_{s}}{\beta''_{z}} = +5.385.$$

$$d_{s} = \frac{-2 d_{s} \beta'_{s} - d_{s} \beta_{s}}{\beta''_{i}} = \frac{-1,109,880 E'}{\beta''_{i}}$$

Therefore,

$$M_2 = \frac{1}{-1.109.880 \ E'} \left\{ \begin{array}{l} \left(\begin{array}{c} -239,198,375 \ E' = 210,013,191 \ E' \\ -101,703,032 \ E' \\ -311,716,223 \ E' \end{array} \right) & 5.385 \\ -311,716,223 \ E' \end{array} \right\} \\ \left(\begin{array}{c} -84,584,500 \ E' - 69,224,626 \ E' - 239,198,375 E' \\ -202,036,278 \ E' \\ -262,260,904 \ E' \\ -262,260,904 \ E' \\ -37,029,904 \ E' \\ -4,881,977 \ E' \end{array} \right) & (-1.524) \\ \left(\begin{array}{c} -122,469,568 \ E + 22,147,927 \ E' - 36,250,500 \ E' \\ -27,029,904 \ E' \\ -2,586,360 \ E' \end{array} \right) \\ \left(\begin{array}{c} -5,385 \ (2.33) \\ -2,596,360 \ E' \end{array} \right) & (-131,890,304 \ E') \end{array} \right)$$

$$M_{e} = \frac{1}{1,109,880} \begin{cases} -2,966,675,110 \\ +893,130,719 \\ -163,602,045 \end{cases} + \frac{12.565}{2,590,360} (-131,890,304)$$

$$M_z = -2015.66 - 639.68 = -2656.$$

The moments, produced by the loads in the respective spans, can be easily deduced, as follows:

First span loaded-

$$M_{s} = \frac{-d_{s} \beta_{s}^{"} \beta_{s}^{"}}{(c_{s} \beta_{s}) \beta_{s} \beta_{s}^{"}} B_{s} + \frac{-1}{(d_{s} \beta_{s}^{"})} X_{s}^{"} d_{s}.$$

$$Or, \quad \begin{cases} +\frac{1}{1,109,880} (-101,703,032) (5,385) \\ +\frac{12,565}{2,590,360} (-131,890,304) \end{cases} = \begin{bmatrix} -494 \\ -640 \\ \hline -1134 \end{bmatrix}$$

$$M_{2} = \frac{-1}{(d_{5} \beta_{1}^{"})} \left\{ \begin{array}{c} (A_{2} + X_{2}^{'}) \ d_{4} \\ (X_{2}^{"} + B_{2}) \ d_{3} \end{array} \right\}$$

$$M_{2} = \frac{1}{1,109,880} \left\{ \begin{array}{c} (-239,198,375 - 210,013,191) \ (5,385) \\ (-202,036,278 - 239,198,375) \ (-1,524) \end{array} \right\} = \dots -157$$

Third span loaded-

$$M_{2} = \frac{1}{(d_{5} \beta_{1}'')} \left\{ \begin{pmatrix} (A_{s} + X_{s}') \ d_{s} \\ (X_{s}'' + B_{s}) \end{pmatrix} \right\}$$

$$M_{2} = \frac{1}{1,109,880} \left\{ \begin{pmatrix} -84,584,500 - 60,224,626) & (-1,524) \\ (-27,029,904 - 36,250,500) \end{pmatrix} \right\} = \dots + 142$$

Fourth span loaded-

$$M_{s} = \frac{-1}{(d_{s} \beta_{i}^{\prime})} \left\{ A_{s} + X_{s}^{\prime} \right\}$$

$$M_{s} = \frac{1}{1,109,880} \left\{ -\frac{122,469,568}{+22,147,927} \right\} = \frac{-122,469,568}{2000} = \frac{-1$$

THE SAME EXAMPLE WITH I CONSTANT.

$$\begin{split} M_2 &= \frac{-c_z}{d_s \ l_i} \begin{cases} A_z \ d_s + B_z \ d_s \\ A_s \ d_s + B_s \end{cases} \right\} \ + \ \frac{-d_s}{c_s \ l_s} \ B_t \\ c_i &= 0, \ c_z = 1, \ c_s = -3.6, \ c_s = +13.4, \ c_6 = -\frac{2901.6}{52} \\ d_i &= 0, \ d_z = 1, \ d_s = -3.6, \ d_s = +13.4, \ d_s = -\frac{2901.6}{52} \\ A_2 &= -\frac{1}{4} \ w_z \ l_z^3 \qquad A_s = -\frac{1}{4} \ v_s \ l_s^3 \qquad A_s = -\frac{1}{4} \ w_s \ l_s^3 \\ B_i &= -\frac{1}{4} \ w_i \ l_i^3 \qquad B_z = -\frac{1}{4} \ w_z \ l_z^3 \qquad B_s = -\frac{1}{4} \ v_s \ l_s^3 \end{split}$$
Then,
$$M_2 &= \frac{1}{2901.6} \begin{cases} -\frac{1}{4} \ w_s \ l_z^3 \ (-3.6) - \frac{1}{4} \ v_s \ l_s^3 \end{cases}$$

 $-\frac{1}{l} w_{i} l_{i}^{s} - \frac{1}{l} w_{i} l_{i}^{s}$ (13.4)

$$M_{2} = \frac{1}{2901.6} \left\{ \begin{array}{l} -919,993.7 \ w + 247,162.5 \ w \\ + 247,162.5 \ v - 68,656.2 \ v \\ - 35,152,0 \ w - 471,036.8 \ w \end{array} \right\} = -2587$$

The partial moments are easily obtained, as follows:

Second span loaded-

$$M_2 = \frac{1}{2901.6} (-919,993.7 \ w + 247,162.5 \ w) = . . . -1554$$

Third span loaded-

$$M_2 = \frac{1}{2901.6} (247,162.5 \ v - 68,656.2 \ v) = \dots + 136$$

Fourth span loaded-

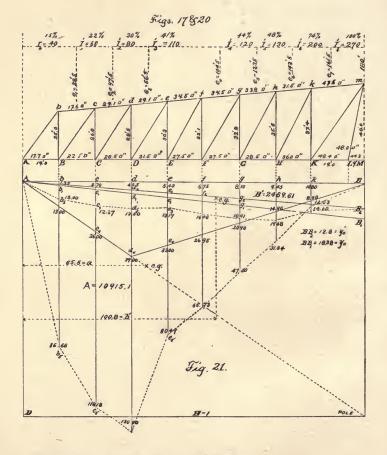
Fourth span loaded—
$$M_2 = \frac{1}{2901.6} (-35{,}126.0 w) = \dots \dots - \frac{81}{2587}$$

$$Sum = -2587$$

Spans loaded.	Variable I .	Constant I .	Difference.	Per cent. of Variable I Results.
FIRST.	— 1134	— 1088	46	.041
- SECOND.	— 1574	— 1554	20	.012
THIRD.	+ 142	+ 136	+ 6	.043
Fourth.	- 90	- 81	9	.100
ALL.	— 2656	2587	69	.026

EXAMPLE 3.—THE SABULA DRAW.

We now propose to compare the bending moments in the Sabula Draw, considering the moment of inertia as constant and variable, respectively.



Variable moment of inertia-

From Fig. 20, is obtained the cross section of each mem-

ber, and the length of each vertical; the moments of inertia for each section are readily deduced as follows: neglecting the moment of inertia of the section of the member about its own axis, it being very small, in comparison with the moment of inertia of the truss section.

FIRST SPAN.

$$I_{o} = (17.7 + 17.6) (12.5)^{2} \div 144 = 40 = I_{l}^{2}$$

$$I_{l} = (29.1 + 22.5) (13.4)^{2} \div 144 = 60 = I_{d}^{2}$$

$$I_{l} = (29.1 + 28.5) (14.25)^{2} \div 144 = 80 = I_{d}^{2}$$

$$I_{d} = (34.5 + 31.5) (15.15)^{2} \div 144 = 110 = I_{d}^{2}$$

$$I_{d} = (33.8 + 27.5) (16.05)^{2} \div 144 = 120 = I_{d}^{2}$$

$$I_{d} = (31.5 + 28.5) (17.8)^{2} \div 144 = 130 = I_{d}^{2}$$

$$I_{d} = (47.5 + 36.0) (18.7)^{2} \div 144 = 200 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 + 270 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 + 270 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 + 270 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 = 130 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 = 200 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 = 130 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 = 130 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 = 130 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 = 130 = I_{d}^{2}$$

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$$I_{l} = (47.5 + 48.0) (20.0)^{2} \div 144 = 130 = I_{d}^{2}$$

$$I_{l} = (47.5 + 48.0) (20$$

From (124), we have,

$$M_2 = \frac{1}{2 \beta_2'} (A_2 + B_1 + X_2' + X_1'')$$

In which,

$$\begin{split} X_{2}' = & H_{2} \; \theta_{1} - F_{2} \; \Sigma \; P_{2} \; l_{2} \; (1 - k_{2}) \; \theta_{1} \\ X_{1}'' = & - H_{1}' \; \theta_{2} - 2 \; F_{1}'' \; \Sigma \; P_{1} \; l_{1} \; (1 - k_{1}) \; \theta_{2} \\ v = & 1 \\ H_{2} = & \sum_{v} \; \left\{ \frac{\Sigma \; P_{2} \; (c_{v} - a_{2})^{3}}{l_{2}} \; + \; \frac{3 \; (l_{2} - e^{-})}{l_{2}} - \Sigma \; P_{2} \; (c_{v} - a_{2})^{2} \right\} \\ v = & l_{2} \end{split}$$

$$H_{i} = \sum_{v=l_{s}} \int_{c}^{t} \left\{ \frac{\sum P_{i} (e_{v} - a_{i})^{3}}{l_{i}} - \frac{3 e_{v}}{l_{i}} \sum P_{i} (e_{v} - a_{i})^{2} \right\}$$

$$v = l_{s}$$

$$E_{i} = E I_{i} (l_{i} + F_{i}^{"}) + 6 E I_{i} (l_{2} + F_{2})$$

$$v = 1$$

$$F_{i} = \sum_{v=l_{s}} \int_{c}^{t} \frac{e_{v}^{3}}{l_{i}^{2}}$$

$$v = l_{s}$$

$$v = l_{s}$$

$$v = l_{s}$$

$$F_z = \begin{array}{c} v = 1 \\ \Sigma & \stackrel{z}{\bigtriangleup} \left\{ \begin{array}{c} l_z - (l_z - e_v)^3 \\ l_z \end{array} \right\} \\ v = l_z \end{array}$$

$$\overset{\circ}{\triangle}_{r} = \overset{\mathring{I}_{l}}{\overset{\circ}{I}_{r-l}} - \overset{\mathring{I}_{l}}{\overset{\circ}{I}_{r}} \qquad \overset{\acute{\triangle}}{\triangle}_{r} = \overset{\acute{I}_{l}}{\overset{\acute{I}_{l}}{I}_{r-l}} - \overset{\acute{I}_{l}}{\overset{\acute{I}_{l}}{I}_{r}}$$

We will first compute those co-efficients that do not depend upon the loading. For a uniform load over all,

$$H_{z} = \frac{v = 1}{\sum_{r}^{2} \bigwedge_{r}^{2} \frac{w_{z} c_{r}^{3}}{4} \left\{ 4 - 3 \frac{e_{r}}{l_{z}} \right\}}$$

$$v = l_{z}$$

$$H_{i} = \sum_{v=l_{i}}^{v=1} \left(\frac{3w_{i}}{4} \left(-\frac{e_{r}^{4}}{l_{i}} \right) \right)$$

First span-

Multiplying this sum by $\frac{3 w_i}{4 l_i}$, we obtain $-2,615,286 w_i = H_i'$.

Second span-

Multiplying this sum by $\frac{w_2}{4}$, we obtain $-723,449 w_2 = H_2$.

First span-l₂=180'-

 $l_i^2 = 32,400$) 4,580,143 (141.36= F_i^* .

Second span- $l_2=180$.

Next find the values of X_2 and X_1 , for a uniform load over all.

$$\begin{array}{lll} \theta_{j}=270 & E', \ \theta_{z}=40 & E', \ \text{in which} \ E'=6 & E. \\ + \ H_{z} \ \theta_{i}=-723,449 \ (270) \ E' \ w_{z}=-195,331,230 \ w_{z} \ E' \\ -\frac{F_{z}^{''}}{2} \ w_{z} \ l_{z}^{2} \ \theta_{i}=+\frac{103.33}{2} \\ & (32,400) \ (270) \ E' \ w_{z}=+451,965,420 \ w_{z} \ E' \\ -F_{i}^{''}\theta_{z} w_{i} \ l_{z}^{''}=-141.36 \ (32,400) \ (40) w_{i} \ E'=-183,202,560 \ w_{i} \ E' \\ -H_{i}^{''}\theta_{z}=+2,615,286 \ (40) \ w_{i} \ E'=-194,611,440 \ w_{i} \ E' \\ -78,591,120 \ w_{i} \ E' \\ -78,591,120 \ w_{i} \ E' \\ A_{z}=-\frac{1}{4} \ w_{z} \ l_{z}^{2} \ 270 \ E'=-393,660,000 \ E' \ w_{z}. \\ B_{i}=-\frac{1}{4} \ w_{z} \ l_{z}^{3} \ 40 \ E'=-58,320,000 \ E' \ w_{z}. \\ Then, \\ M_{z}=\frac{1}{51456E'} \begin{cases} -393,660,000 \ w_{z} \ E' \\ -58,320,000 \ w_{z} \ E' \\ -78,591,120 \ w_{z} \ E' \\ +256,634,190 \ w_{z} \ E' \end{cases} = (\text{if } w_{i}=w_{z}=w)-5322 \ w. \end{cases}$$

First span alone loaded-

$$M_{2} = \frac{1}{51456} \left\{ \frac{-\begin{array}{c} 58,320,000 \ w_{i} \\ -\begin{array}{c} 78,591,120 \ w_{i} \\ \hline -\begin{array}{c} 137,000,000 \ w_{i} \end{array} \right\} = -2661 \ w_{i}.$$

Second span alone loaded-

$$M_s = rac{1}{51456} iggl\{ rac{-393,660,000}{+256,834,190} rac{w_s}{w_z} iggr\} = -2661 \ w_s.$$

EXAMPLE 3a.

This is the same as Example 3, with the moment of inertia considered as constant.

From (142),

$$M_{2} = \frac{A_{2} + B_{i}}{2(l_{i} + l_{2})} = \frac{A_{2} + B_{i}}{720}$$

$$A_{2} = -\frac{1}{4} w_{2} l_{2}^{3} = -145,800 w_{2}.$$

$$B_{i} = -\frac{1}{4} w_{i} l_{i}^{3} = -145,800 w_{2}.$$

Therefore,

$$M_2 = \left\{ \begin{array}{l} -2025 \ w_i \\ -2025 \ w_i \end{array} \right\}$$
 and if $w_i = w_2 = w$, $M_2 = -4050 \ w$.

No. of Loaded Span.		M_2 Bending Moment Constant I.	Difference.
FIRST.	— 2661 w,	- 2025 w,	$636 w_i$
SECOND.	$-2661 w_{z}$	- 2025 w ₂	$636 w_2$
Вотн.	— 5322 w	- 4050 w	1272 w

EXAMPLE 4.—CONCENTRATED LOADS.

Let us consider the Sabula Draw again, but with single concentrated loads, instead of uniform loads. Supports level. $a_i=57$. $a_2=123$.

See pages 73 and 74 for the value of M_2 .

First span— $l_1 = 180'$. $a_2 = 57'$.

From Example 3:

$$\beta'_{z}$$
=40 (180+141.36) E' +270 (180-132.32) E' =+25,728 E' .
 F''_{z} =+141.36. F'_{z} =103.33. F_{z} =-132.32.
 θ_{z} =270 E' . θ_{z} =40 E' .

The values of H_2 and H_i' depend upon the position of the loads. Let us take a load in the first span 57' from the left support, and one in the second span 123' from the left support; then we have for H_i' and H_2 , the following:

$$v=1 \\ H_i' = \sum_{v} \triangle_v^i \left\{ \frac{\sum_i P_i \left(e_v - a_i \right)^2}{l_i} - \frac{3 e_v}{l_i} \stackrel{\Sigma}{=} P_i \left(e_v - a_i \right)^2 \right\} \\ v = l_i \\ e_i = 28.5 \quad e - a_i. \\ e_2 = 47.5 \\ e_3 = 66.5 \quad + \quad 9.5 \quad \triangle_s = + \quad 0.921 \quad - \quad 88 P_i \\ e_4 = 104.5 \quad + \quad 47.5 \quad \triangle_4 = + \quad 0.204 \quad - \quad 680 P_i \\ e_5 = 123.5 \quad + \quad 66.5 \quad \triangle_5 = + \quad 0.173 \quad \cdot - \quad 1,292 P_i \\ e_6 = 142.5 \quad + \quad 85.5 \quad \triangle_6 = + \quad 0.727 \quad - \quad 10,097 P_i \\ e_7 = 161.5 \quad + \quad 104.5 \quad \triangle_7 = + \quad 0.350 \quad - \quad 8,068 P_i \\ \end{cases}$$

Note. $-e_c$ must always be greater than or equal to a_c .

 $-20225 P_{i}=H'_{i}$.

Second span— $l_2=180'$. $a_2=123'$.

$$H_z = rac{v = 1}{\sum igsim_v^2} \left\{ rac{\Sigma P_z \left(e_v - a_z
ight)^3}{l_z} + rac{\Im \left(l_z - e_v
ight)}{l_z} \, \, \Sigma \, P_z \, \left(e_v - a_z
ight)^2
ight\}$$

$$e_{1} = 18.5$$
 $e_{2} = 37.5$
 $e_{3} = 56.5$
 $e_{4} = 75.5$
 $e_{5} = 113.5$
 $e_{6} = 132.5$
 $e_{7} = 151.5$
 $e_{8} = 28.5$
 $e_{1} = -0.333$
 $e_{8} = -0.333$

Note.— e_v must always be greater than or equal to a_2 .

$$X_{s}' = \begin{cases} + H_{s} \theta_{r} = -185 (270) E' P_{s} = -49,950 E' P_{s} \\ - F_{s}' P_{s} (l_{s} - a_{s}) \theta_{r} = \\ + 103,33 (57) (270) E' P_{s} = +1,590,248 E' P_{s} \\ + 1,540,298 E' P_{s} \end{cases}$$

$$X_{t}'' = \begin{cases} - H_{t}' \theta_{s} = +20,225 P_{t} (40) E' = +809,000 E' P_{t} \\ -2 F_{t}'' P_{t} (l_{t} - a_{t}) \theta_{s} = \\ -141.36 (123) (80) E' P_{t} = -1,390,928 E' P_{t} \\ -581,928 E' P_{t} \end{cases}$$

From (i) and (j), $A_z = -P_z l_z^2 (2 k - 3 k^2 + k^3) 6 E' I_i$, Or, by Table I,

$$A_2 = -P_2 l_z^2 (0.285,144) 270 E' = -2,494,269 E' P_2.$$

 $B_i = -P_i l_i^2 (0.285,144) 40 E' = -369,520 E' P_i.$

Therefore,

$$M_{2} = \frac{1}{51456 E'} \begin{cases} -2,494,269 E' P_{2} \\ +1,540,298 E' P_{2} \\ -369,520 E' P_{1} \\ -581,928 E' P_{1} \end{cases} = \begin{cases} -185 P_{2} \\ -18.5 P_{1} \end{cases}$$
Or, if $P_{1} = P_{2}$,
$$M_{2} = -37.0 P$$
.

First span alone loaded—

$$M_2 = rac{1}{2 eta_2'} \left(B_i + X_i''
ight)$$

Or,

$$M_2 = \frac{1}{51456} \left\{ -951,448 \ P_i \right\} = 18.5 \ P_i.$$

Second span alone loaded--

$$M_{z}=rac{1}{2\leftert eta_{z}^{\prime}}\left\{ A_{z}+X_{z}^{\prime}
ight.
ight\}$$

Or,

$$M_z = \frac{1}{51456} \left\{ -953,961 \ P_z \right\} = 18.5 \ P_z.$$

Since P_1 and P_2 are symmetrical about the center of the draw, the moments over the center pier should be equal, as shown above. If our work was absolutely correct, and decimals had been used to several places, the quantities in the parentheses above would have been the same; as it is, the quotients are practically equal.

* EXAMPLE 4 BY GRAPHICS.

For any load in the first span of a two span girder, we have, from *Greene's Trusses and Arches*, Part II,

$$\Sigma \frac{x_i \ y_i}{E \ \dot{I}} - \Sigma \frac{x_i' \ y_i'}{E \ \dot{I}} = \Sigma \frac{x_2' \ y_2'}{E \ \dot{I}}$$

In which:

 x_i =The distance from the left support to the ordinate y_i of the equilibrium polygon $A \subset B$. Fig. 21, page 72.

 y_i =Any ordinate of the polygon $A \ C \ B$.

 x_i =The distance from the left support to the ordinate y_i of the polygon $A B B_i$.

 y_i =Any ordinate of the polygon $A B B_i$.

 x_i =The distance from the *right* support to any ordinate y_i of a polygon in the second span, similar to ABB_i .

^{*} For this excellent graphical method, the author is indebted to R. H. Brown, C. E., First Assistant Engineer of Boston Bridge Works.

 y_2' =Any ordinate of the above polygon.

I=The moment of inertia of the cross section of the girder at any ordinate y'_i .

 $\stackrel{?}{I}$ =The moment of inertia of the cross section of the girder at any ordinate y_2 .

E=The modulus of elasticity.

Since $l_i = l_2 = l$, and I = I = l, we have, considering E as constant,

$$\Sigma \frac{x_i \ y_i}{I} = 2 \ \Sigma \frac{x_i' \ y_i'}{I}$$

Now, $\frac{x_i y_i}{I}$ equals an area multiplied by the distance of its center of gravity from the left support, and hence we may write $\frac{x_i}{I} = A a$, in which A represents the area and a the c. g. distance.

In like manner, $2 \ge \frac{x_i' y_i'}{I} = 2 B b$. Therefore, we may write, A = 2 B b.

With any scale, lay off the horizontal line AB, Fig. 21, and divide it into panel points at b, c, d, etc.

With any scale, preferably a large scale, lay off a load line \boldsymbol{A} \boldsymbol{D} , equal unity, and assume \boldsymbol{H} also equal unity, then, assuming the pole in such a position that its closing line will be horizontal, construct the equilibrium polygon \boldsymbol{A} \boldsymbol{C} \boldsymbol{B} and scale its ordinates (the lengths are given in Fig. 21.)

Not knowing the proper position of the closing line, let us assume a position as A B_2 , making B $B_2=12.8$, and scale the ordinates b b_2 , c c_2 , etc.

Divide each ordinate of the polygon A C B by its proper I (given in per cent. of the center I above Fig. 21), and lay

off the results downward from AB, forming the polygon $Ab_5c_5...B$. The ordinates are 86.66, 118.18, etc., as shown in Figure 21, page 72.

Now, find the area of this figure, either by computation or the planimeter, and also its center of gravity. The center of gravity is readily found by cutting the polygon out of stiff card board and balancing it upon a needle point.

The area =
$$10915.1 = A$$
, and $a = 65.5$.

Proceed in like manner with the polygon $A B B_2$, deducing the figure $A b_3 c_3 \dots B_2$, which has an area of 2469.61=B', and b'=100.8.

Let B B_i represent the true magnitude of the pier ordinate y_0 , then from the triangles A B B_2 and A B B_i , we have,

$$y_o': y_o :: e \ e_z := e_i$$
, or, $e \ e_i = \frac{y_o}{y_o'} \ e \ e_z$. Then,
$$\Sigma \frac{x_i' \ (e \ e_i)}{I} = \frac{y_o}{y_o'} \ \Sigma \frac{x_z' \ (e \ e_z)}{I}$$
, or, $B \ b = \frac{y_o}{y_o'} \ B' \ b'$.
Hence, $A \ a = 2 \frac{y_o}{y_o'} \ B' \ b'$, and $y_o = \frac{(A \ a) \ y_o'}{2 B' \ b'}$,

Which becomes

$$y_0 = \frac{10915.1 \times 65.5 \times 12.8}{2 \times 2469.61 \times 100.8} = 18.38.$$

Since H=1 and $P_i^s=1$, the bending moment over the center pier is $-18.38 P_i^s$. We obtained, by computation, $-18.5 P_i^s$, which shows the graphical method to be accurate enough for all practical purposes.

It would have been more correct to have taken the ordinates y at the points where the values of I change, but the result would have been but little different.

The above computations can be very readily made by means of Thatcher's Calculating Instrument.

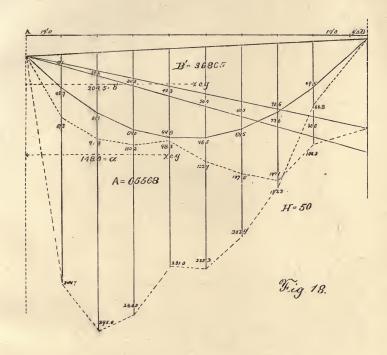


Fig. 17, page 72, and Fig. 18 are diagrams for the Sabula Draw, with a concentration at each apex in both spans.

$$(P_i^i = P_i^z = P_z^i = P_z^z = &c. = P = 23,000 \text{ lbs.})$$

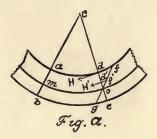
$$A = B b$$
, or $A = \frac{y_0}{y_0'} B' b$. Therefore,

$$y_0 = \frac{(A \ a) \ y_0'}{B' \ b'} = \frac{65568 \times 148 \times 94.6}{36865 \times 2045} = 121.800.$$

Since H=50, the bending moment over the pier is $121,800\times50=\ldots$ 6,090,000



APPENDIX.



* The general equation of the elastic line can be deduced as follows:

Vertical forces acting upon a girder cause a change of shape, lengthening the originally parallel fibres on one side, and shortening or compressing them on the other. Between the lengthened and shortened fibres there is a plane which undergoes no change in length; the centre line of this plane is called the *neutral axis or the elastic line*.

Thus, in Fig. (a), m o is the neutral axis, the fibres above being compressed, and those below lengthened. Upon the three following hypotheses we shall deduce the equation of the *elastic line*.

- I. All planes perpendicular to the axis before the bending or flexure, preserve, during the bending, their perpendicularity and their forms as planes.
- II. The change in length of a body subjected to a force, is, within certain limits, called the elastic limits, proportional to the intensity of the force.

^{*} See Merriman's Theory and Calculation of Continuous Beams.

III. The change of shape is so small that the length of the neutral axis is sensibly the same as its horizontal projection.

In Fig. (a), we have a longitudinal section of a portion of a bent beam; the two planes a b and d e origininally parallel, remaining perpendicular to the neutral axis or elastic line m o, and intersecting in e the centre of curvature. Hence, drawing f g parallel to g g through g, the lines g g g g e, etc., denote the elongation of the fibres, and we see, from the figure, that

o
$$d:o \ d':: d \ f:d' \ f' \ \dots \ \dots \ \dots \ \dots \ (22a)$$

or the change of length in the fibres is proportional to their distance from the neutral axis. This is a consequence of the first hypothesis.

Designating by H and H' the force acting in the fibres df and d'f', the second hypothesis gives

$$H:H'::df:d'f'\ldots\ldots\ldots\ldots\ldots\ldots\ldots$$
 (22b)

Combining (22a) and (22b), we have

$$H:H'::o\ d:o\ d'$$
 (22c)

or, the horizontal forces are directly proportional to their distances from the neutral axis.

Denote the distance of any fibre from the neutral axis by z, the stress in it by H', the distance of the remotest fibre by e, and its stress by H; then, from (22c), we obtain

Thus far the cross-section of the fibres has been considered as unity. If the actual area is a, the force is $\frac{Haz}{e}$. Each of these forces H' tend to turn the beam around o with a lever arm o d' or z, hence the moment of the force is Haz

 $\frac{H~a~z}{e} imes z = \frac{H~a~z^2}{e}$, and the sum of all the moments is

$$M_x = \frac{H}{e} \Sigma a z^2 \dots (22e)$$

 M_x meaning the bending moment at any section x_r of the beam in the span l_r .

Since $\Sigma a z^2$ is the expression for the moment of inertia of the section a b, (22*e*) becomes

or the moment of the internal forces equals the stress in the remotest fibre times the moment of inertia of the section divided by the distance of the remotest fibre from the neutral axis.

The line d f denotes the change of length in the fire a d, due to the force H, hence, if E be the co-efficient of elasticity,

Designating the radius c o by γ_r , we have, from similar figures, o d f and c a d. (m o = a d).

Substituting this value of e in (22f), we have

The radius of curvature of any plane curve, whose length is u_r and co ordinates x_r and y_r , is

$$\gamma_r = \frac{d \ u_r^3}{d \ x_r \ d^2 \ y_r} . \qquad (22l)$$

According to the third hypothesis, $du_r = dx_r$, and (221) becomes

$$\gamma_r = \frac{d \ x_r^2}{d^2 \ y_r} \quad \dots \quad (22m)$$

Substituting (22m) in (22k), it becomes

$$\frac{d^{z}}{d} \frac{y_{r}}{x_{r}^{z}} = \frac{M_{x}}{E} \frac{1}{I_{x}} \dots (22)$$

Which is the differential equation of the elastic line, applicable to all bodies subjected to flexure which fulfill the conditions imposed by the *third hypothesis*. The values of *E* and *I* may be different for each and every section.

If (22) be integrated, it becomes

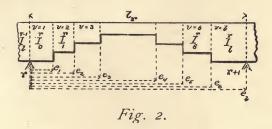
If $x_r = 0$, then $C = \frac{d y_r}{d x_r} = t_r$, and we have

Integrating again, (24) becomes

$$y_r = C' + t_r \int_0^{x_r} d x_r + \frac{1}{E} \int_0^{x_r} \frac{M_x d x_r}{I_x} \int_0^{x_r} d x_r \dots$$
 (25)

If $x_r = 0$, then $C' = h_r$, and we have

$$y_r = h_r + t_r x_r + \frac{1}{E} \int_0^{x_r} \frac{M_x}{I_x} \frac{d}{dx_r} x_r \int_0^{x_r} dx_r \dots \dots (26)$$



By examining Fig. 2, which represents a continuous girder having a variable cross-section, and consequently a variable moment of inertia, the following integration will be clearly understood:

* Integration of
$$\int_{0}^{x_r} \frac{r}{M_x} \frac{d}{dx_r} \frac{x_r}{I_x}$$

$$\int_{0}^{\infty} \frac{M_{x}^{r} d x_{r}}{I_{x}} = \frac{1}{I_{0}} \int_{0}^{e_{t}} M_{x}^{r} d x_{r} + \frac{1}{I_{t}} \int_{e_{t}}^{e_{z}} M_{x}^{r} d x_{r} . . .$$

$$\ldots + \frac{1}{I_x} \int_{e_x}^x M_x \ dx_r \ldots (27)$$

But,
$$\frac{1}{\tilde{I}_{v}} \int_{v}^{e_{v+t}} M_{x}^{r} d x_{r} = \frac{1}{\tilde{I}_{v}} \int_{o}^{e_{v+t}} M_{x}^{r} d x_{r} - \frac{1}{\tilde{I}_{v}} \int_{o}^{e_{v}} M_{x}^{r} d x_{r} (28)$$

Therefore, we can write for (27),

^{*} See Weyrauch's Continuirlichen und Einfachen Träger, p. 168.

$$\int_{0}^{T} \frac{1}{I_{x}} M_{x}^{r} d x_{r} = \frac{1}{I_{0}} \int_{0}^{e_{f}} M_{x}^{r} d x_{r} + \frac{1}{I_{1}} \int_{0}^{e_{2}} M_{x}^{r} d x_{r} ... + \frac{1}{I_{x-1}} \int_{0}^{e_{x}} M_{x}^{r} d x_{r}$$

$$- \frac{1}{I_{1}} \int_{0}^{e_{f}} M_{x}^{r} d x_{r} - \frac{1}{I_{2}} \int_{0}^{e_{2}} M_{x}^{r} d x_{r} ... - \frac{1}{I_{x}} \int_{0}^{e_{x}} M_{x}^{r} d x_{r}$$

$$+ \frac{1}{I_{x}} \int_{0}^{e_{x}} M_{x}^{r} d x_{r} (29)$$

Which reduces to

$$\int_{0}^{x_{r}} \frac{1}{I_{x}^{r}} M_{x}^{r} dx_{r} = \frac{1}{I_{x}^{r}} \int_{0}^{x_{r}} M_{x}^{r} dx_{r}$$

$$v=1$$

$$+ \frac{\Sigma}{v=x_{r}} \left(\frac{1}{I_{v-s}} - \frac{1}{I_{v}^{r}} \right) \int_{0}^{e_{v}} M_{x}^{r} dx_{r} \dots (30)$$

Substituting (30) in (26), it reduces to

$$y_{r} = h_{r} + t_{r} x_{r} + \frac{1}{E} \left\{ \frac{1}{I_{x}} \int_{0}^{x_{r}} dx_{r} \int_{0}^{x_{r}} M_{x} dx_{r} \right.$$

$$v = 1$$

$$+ \frac{\Sigma}{v = x_{r}} \left(\frac{1}{I_{v-f}} - \frac{1}{I_{v}} \right) \int_{0}^{x_{r}} dx_{r} \int_{0}^{x_{r}} M_{x} dx_{r} \right\} . . . (31)$$

From 8, we have $\stackrel{r}{M}_x = M_r + S_r x_r - \Sigma P_r (x_r - a_r) \cdot ... a \leq x \cdot ... (8)$

Hence,

$$\int_{0}^{x_{r}} M_{x} d x_{r} = M_{r} x_{r} + \frac{1}{2} S_{r} x_{r}^{2} - \frac{1}{2} \Sigma P_{r} (x_{r} - a_{r})^{2} \dots a \leq x \dots (32)$$

And,

$$\int_{0}^{x_{r}} d_{x} x_{r} \int_{0}^{x_{r}} M_{x} dx_{r} = \frac{1}{2} M_{r} x_{r}^{2} + \frac{1}{6} S_{r} x_{r}^{2}$$

$$- \frac{1}{6} \Sigma P_{r} + v_{r} - a_{r}^{3} \cdot \cdot \cdot a \leq x \cdot \cdot (33)$$

Now,

$$\int_{0}^{x_{r}} dx_{r} \int_{0}^{e_{v}} M'_{x} dx_{r} = \int_{0}^{e_{v}} dx \int_{0}^{e_{v}} M'_{x} dx,$$

$$+ \int_{e_{v}}^{x_{r}} dx_{r} \int_{0}^{e_{v}} M'_{x} dx_{r} \dots \dots (34)$$

But,

$$\int_{0}^{c_{r}} M_{x}^{r} dx_{r} = M_{r} e_{r} + \frac{1}{2} S_{r} c_{r}^{2} - \frac{1}{2} \Sigma P_{r} (e_{r} - a_{r})^{2} \dots a \leq e \dots (35)$$

Hence,

$$\int_{0}^{e_{v}} dx_{r} \int_{0}^{e_{v}} M_{x}^{r} dx_{r} = \frac{1}{2} M_{r} e_{v}^{2} + \frac{1}{6} S_{r} e_{v}^{5} - \frac{1}{6} \Sigma P_{r} (e_{v} - a_{r})^{3} \dots a \leq e \dots (36)$$

$$\int_{e_{r}}^{x_{r}} dx_{r} = x_{r} - e_{v},$$

Therefore,

$$\int_{e_{v}}^{x_{r}} \int_{0}^{e_{v}} M_{x}^{r} dx_{r} = \frac{1}{2} (x_{r} - e_{v}) \left\{ 2 M_{r} e_{v} + S_{r} e_{v}^{2} - \sum P_{r} (e_{r} - a_{r})^{2} \right\} \dots (37)$$

Substituting (33), (34), (36) and (37) in (31), it becomes $y_{r} = h_{r} + t_{r} x_{r} + \frac{1}{6 E \tilde{I}_{x}} \left\{ 3 M_{r} x_{r}^{2} + S_{r} x_{r}^{3} - \Sigma P_{r} (x_{r} - a_{r})^{3} \right\}$ $+ \frac{1}{6 E \tilde{I}_{x}} \sum_{v = x_{r}} \left(\frac{\tilde{I}_{x}}{\tilde{I}_{v-i}} - \frac{\tilde{I}_{x}}{\tilde{I}_{v}} \right) \left\{ 3 M_{r} e_{v}^{2} + S_{r} e_{v}^{3} - \Sigma P_{r} (e_{v} - a_{r})^{3} \right\}$

$$+3(x_r-e_v)\left[2M_re_v+S_re_v^2-\Sigma P_r(e_v-a_r)^2\right]_{v=0}^{k}$$
 . . . (38)

Which reduces to

$$y_{r} = h_{r} + t_{r} x_{r} + \frac{1}{6 E I_{x}^{r}} \left\{ \beta M_{r} x_{r}^{2} + S_{r} x_{r}^{3} - \Sigma P_{r} (x_{r} - a_{r})^{3} \right\}$$

$$+ \frac{1}{6 E I_{x}^{r}} \sum_{v=-x_{r}}^{y=1} \left(\frac{I_{x}^{r}}{I_{v-1}^{r}} - \frac{I_{x}^{r}}{I_{v}} \right) \left\{ \beta M_{r} e_{v} (2 x_{r} - e_{v}) + S_{r} e_{v}^{2} (\beta x_{r} - e_{v}) - \Sigma P_{r} (e_{v} - a_{r})^{3} - \beta (x_{r} - e_{v}) \Sigma P_{r} (e_{v} - a_{r})^{2} \right\} (39)$$

If we make $x_r = l_r$, then $y_r = h_{r+l}$, $e_x - e_l$, and $\vec{I_x} = \vec{I_l}$.

Let
$$\frac{I_t^r}{I_{v-t}^r} - \frac{I_t^r}{I_v^r} = \stackrel{r}{\triangle}_v$$
, then, from (39), we have

$$\begin{split} h_{r+t} &= h_r + t_r \ l_r + \frac{1}{6 \ E \ I_t} \ \left\{ \beta \ M_r \ l_r^2 + S_r \ l_r^3 - \Sigma \ P_r \ (l_r - a_r)^3 \right\} \\ v &= 1 \end{split}$$

$$+ \frac{1}{6 L I_{l}} \sum_{v=l_{r}}^{r} \left\{ 3 M_{r} e_{v} \left(2 l_{r} - e_{v} \right) + S_{r} e_{v}^{2} \left(3 l_{r} - 2 e_{v} \right) \right.$$

$$- \Sigma P_r (e_v - a_r)^s - \beta (l_r - e_v) \Sigma P_r (e_v - a_r)^{\frac{1}{2}} . . . (40)$$
From (10),

But, since $k_r = \frac{a_r}{l_r}$, this becomes

Substituting (10a) in (40), and solving for t_r , we obtain

$$t_{r} = \frac{h_{r+s} - h_{r}}{l_{r}} - \frac{1}{6 E I_{t} l_{r}} \left\{ 2 M_{r} l_{r}^{2} + M_{r+s} l_{r}^{2} + \frac{\Sigma P_{r} a_{r} (l_{r} - a_{r}) (2 l_{r} - a_{r})}{+ \Sigma P_{r} a_{r} (l_{r} - a_{r}) (2 l_{r} - a_{r})} \right\}$$

$$- \frac{v = 1}{6 E I_{t} l_{r}} \frac{\Sigma \triangle_{v}}{v = l_{r}} \left\{ 2 M_{r} e_{v} (3 l_{r} - 3 e_{v} - \frac{e_{r}^{2}}{l_{r}}) + M_{r+s} e_{v}^{2} (3 - \frac{2 e_{v}}{l_{r}}) \right\}$$

$$+ e_{r}^{3} \left(\beta - \frac{2 e_{r}}{l_{r}} \right) \Sigma P_{r} \left(l_{r} - a_{r} \right) - \Sigma P_{r} \left(e_{r} - a_{r} \right)^{3} \\ - \beta \left(l_{r} - e_{r} \right) \Sigma P_{r} \left(e_{r} - a_{r} \right)^{2} \right\} (41)$$

Since $a_r = k_r \ l_r$, $a_r \ (l_r - a_r) \ (2 \ l_r - a_r) = l_r^3 \ (2 \ k_r - 3 \ k_r^2 + k_r^3)$, $l_r - a_r = l_r \ (1 - k_r)$, and (41) reduces to

$$t_{r} = \frac{h_{r+s} - h_{r}}{l_{r}} - \frac{1}{6 E I_{t}} \left\{ 2 M_{r} l_{r} + M_{r+s} l_{r} + \sum P_{r} l_{r}^{2} (2 k_{r} - 3 k_{r}^{2} + k_{r}^{3}) \right\}$$

$$v=l$$

$$-\frac{1}{6 E \stackrel{r}{I_{t}} \stackrel{v}{l_{r}}} \stackrel{\Sigma}{\sim} \stackrel{r}{\bigtriangleup}_{v} \left\{ \stackrel{?}{\sim} M_{r} \stackrel{e_{v}}{e_{v}} \left(\stackrel{?}{\sim} l_{r} - \stackrel{?}{\sim} e_{v} + \frac{e_{v}^{2}}{l_{r}} \right) + M_{r+r} \stackrel{e^{2}}{e_{v}} \left(\stackrel{?}{\sim} - \frac{\stackrel{?}{\sim} e_{v}}{l_{r}} \right) \right\}$$

$$+ e_{v}^{2} \left(3 - \frac{2 e_{v}}{l_{r}} \right) \Sigma P_{r} l_{r} \left(1 - k_{r} \right) - \Sigma P_{r} \left(e_{v} - a_{r} \right)^{3} - 3 \left(l_{r} - e_{v} \right) \Sigma P_{r} \left(e_{r} - a_{r} \right)^{2} \right\} (42)$$

Returning to (24),

Substituting (32) and (35) in (30), and the result in (24), it becomes

$$\frac{d y_{r}}{d x_{r}} = t_{r} + \frac{1}{2 E I_{x}^{r}} \left\{ 2 M_{r} x_{r} + S_{r} x_{r}^{2} - \Sigma P_{r} (x_{r} - a_{r})^{2} \right\}
+ \frac{v = 1}{2 E I_{x}^{r}} \frac{\Sigma}{v = x_{r}} \left(\frac{I_{x}}{I_{r-i}} - \frac{I_{x}}{I_{r}} \right)
\left\{ 2 M_{r} e_{r} + S_{r} e_{v}^{2} - \Sigma P_{r} (e_{r} - a_{r})^{2} \right\} \dots (43)$$

Making $x_r = l_r$, then $\frac{d}{d} \frac{y_r}{x_r} = t_{r+i}$, $e_x = e_t$, and $\stackrel{r}{I}_x = \stackrel{r}{I}_t$, and substituting for S_r its value from (10), t_r from (42), and $k_r l_r$ for a_r in (43), it becomes

$$t_{r+i} = \frac{h_{r+i} - h_r}{l_r} - \frac{1}{6 E I_t} \left\{ 2 M_r l_r + M_{r+i} l_r + P_r l_r^2 \left(2 k_r - 3 k_r^2 + k_r^3 \right) \right\}$$

$$+ P_r l_r^2 \left(2 k_r - 3 k_r^2 + k_r^3 \right) \left\{ 2 M_r e_r \left(3 l_r - 3 e_r + \frac{e_r^2}{l_r} \right) + M_{r+i} e_r^2 \left(3 - \frac{2 e_v}{l_r} \right) \right\}$$

$$+ e_r^2 \left(3 - \frac{2 e_v}{l_r} \right) \Sigma P_r l_r \left(1 - k_r \right) - \Sigma P_r \left(e_v - a_r \right)^3 - 3 \left(l_r - e_r \right) \Sigma P_r \left(e_v - a_r \right)^2 \right\}$$

$$+ \frac{1}{6 E I_t} \left\{ 3 M_r l_r + 3 M_{r+i} l_r + 3 \Sigma P_r l_r^2 \left(k_r - k_r^2 \right) \right\}$$

$$+ \frac{v = 1}{6 E I_t} l_r \Sigma \triangle_v \left\{ 3 M_r e_v l_r \left(2 - \frac{e_v}{l_r} \right) + 3 M_{r+i} e_v^2 + l_r^2 \left(k_r - k_r^2 \right) \right\}$$

$$+ 3 \Sigma P_r \left(1 - k_r \right) e_v^2 l_r - 3 l_r \Sigma P_r \left(e_r - a_r \right)^2 \right\} (44)$$

Which reduces to

$$t_{r+i} = \frac{h_{r+i} - h_r}{l_r} \; + \; \frac{1}{6 \; E \; \overset{r}{I_I}} \left\{ \; M_r \; l_r + 2 \; M_{r+i} \; l_r + \Sigma \; P_r \; l_r^2 \; (k_r - k_r^3) \; \right\}$$

$$+\frac{1}{6} \frac{v=1}{L_{t}} \sum_{l_{r}}^{r} \sum_{v=l_{r}}^{r} \left\{ M_{r} e_{v}^{2} \left(3 - \frac{2 e_{v}}{l_{r}}\right) + 2 M_{r+s} \frac{e_{v}^{3}}{l_{r}} + \frac{2 e_{v}^{3}}{l_{r}} \sum_{r} P_{r} l_{r} (1 - k_{r}) - 3 e_{v} \sum_{r} P_{r} \left(e_{v} - a_{r}\right)^{2} + \sum_{r} P_{r} \left(e_{v} - a_{r}\right)^{3} \right\}$$

$$(45)$$

If we were to suppose loads in the $r-1^n$ span at distances $a_{r-j}=k_{r-j}$ l_{r-j} from the left support r-1, we would find in a similar manner, or by *decreasing* the subscripts of (45) by unity,

$$t_{r} = \frac{h_{r} - h_{r-i}}{l_{r-i}} + \frac{1}{6 E I_{i}} \left\{ M_{r-i} \ l_{r-i} + 2 M_{r} \ l_{r-i} + 2 M_{r}$$

Equating (42) and (46), we obtain

v = l.

$$\begin{cases}
M_{r-i} l_{r-i} + 2 M_r l_{r-i} + \Sigma P_{r-i} l_{r-i}^2 (k_{r-i} - k_{r-i}^3) + \Sigma & \sum_{v} \frac{r^{-i}}{\Delta_v} \epsilon_v^2 \\
v = l_{r-i}
\end{cases}$$

$$v = 1 \qquad v = 1 \qquad v = 1 \qquad v = 1 \qquad v = l_{r-i}$$

$$\left(\frac{3}{l_{r-i}} - \frac{2 e_v}{l_{r-i}^2}\right) M_{r-i} + \frac{\Sigma}{2} \sum_{v} \frac{e_v^3}{l_{r-i}^2} 2 M_r + \frac{\Sigma}{2} \sum_{v} \frac{2 e_v^3}{l_{r-i}^2} v = l_{r-i}$$

$$v = 1 \qquad v = l_{r-i}$$

$$v = 1 \qquad v = l_{r-i}$$

$$v = l_{r$$

$$v=1$$

$$\sum_{r} \bigwedge_{r} \left\{ \frac{\sum P_{r} \left(e_{r}-a_{r}\right)^{s}}{l_{r}} + \frac{3 \left(l_{r}-e_{r}\right)}{l_{r}} \sum P_{r} \left(e_{r}-a_{r}\right)^{s} \right\} = H_{r} ...(52)$$

$$v=l_{r}$$

$$v=1$$

$$\sum_{r} \int_{r}^{r} \left\{ \frac{\sum P_{r} (e_{r}-a_{r})^{s}}{l_{r}} - \frac{\beta e_{v}}{l_{r}} \sum P_{r} (e_{r}-a_{r})^{s} \right\} = H_{r}' \quad . \quad (53)$$

$$v=l_{r}$$

Substituting (48), (49), (50), (51), (52) and (53) in (47), transposing and reducing, we obtain

$$- \theta_{r} \theta_{r-j} \frac{h_{r-j} - h_{r}}{l_{r}} + 2 M_{r} l_{r} \theta_{r-j} + M_{r+j} l_{r} \theta_{r-j} + \Sigma P_{r} l_{r}^{2} (2 k_{r} - 3 k_{r}^{2} + k_{r}^{2}) \theta_{r-j} + F_{r} 2 M_{r} \theta_{r-j} + F_{r}' M_{r+j} \theta_{r-j} + F_{r}' \Sigma P_{r} l_{r} (1 - k_{r}) \theta_{r-j}$$

$$- H_{r} \theta_{r-j} + \frac{h_{r} + h_{r-j}}{l_{r-j}} \theta_{r} \theta_{r-j} + M_{r-j} l_{r-j} \theta_{r} + 2 M_{r} l_{r-j} \theta_{r} + 2 M$$

Which reduces to

$$M_{r-i} (l_{r-i} + F'_{r-i}) \theta_r + 2 M_r \left\{ (l_r + F_r) \theta_{r-i} + (l_{r-i} + F''_{r-i}) \theta_r \right\}$$

$$+ M_{r+i} (l_r + F'_r) \theta_{r-i} = - \theta_r \theta_{r-i} \left\{ \frac{h_r - h_{r-i}}{l_{r-i}} + \frac{h_r - h_{r+i}}{l_r} \right\}$$

$$+ \Sigma P_r l_r^2 (2 k_r - 3 k_r^2 + k_r^3) \theta_{r-i} - F'_r \Sigma P_r l_r (1 - k_r) \theta_{r-i}$$

$$- \Sigma P_{r-i} l_{r-i}^2 (k_{r-i} - k_{r-i}^2) \theta_r - F''_{r-i} 2 \Sigma P_{r-i} l_{r-i} (1 - k_{r-i}) \theta_r$$

$$+ H_r \theta_{r-i} - H'_{r-i} \theta_r \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots (55)$$

$$-\theta_r \theta_{r-i} \left\{ \frac{h_r - h_{r-i}}{l_{r-i}} + \frac{h_r - h_{r+i}}{l_r} \right\} = Y_r \dots \dots (56)$$

$$= F' \stackrel{\Sigma}{\Sigma} P \stackrel{l}{\downarrow} (1-k) \stackrel{\theta}{\downarrow} + H \stackrel{\theta}{\downarrow} = X' \qquad (59)$$

$$-F'_{r} \stackrel{\Sigma}{\Sigma} P_{r} l_{r} (1-k_{r}) \theta_{r-i} + H_{r} \theta_{r-i} = X'_{r} ...$$

$$-2 F''_{r-i} \stackrel{\Sigma}{\Sigma} P_{r-i} l_{r-i} (1-k_{r-i}) \theta_{r} - H'_{r-i} \theta_{r} = X''_{r-i}$$

$$\vdots \qquad (59a)$$

$$\theta_{r-1}(l_r+F_r')=\beta_r''\ldots\ldots\ldots\ldots\ldots\ldots(62)$$

Then we can write

$$M_{r-1} \beta_{r-1} + 2 M_r \beta_r' + M_{r+1} \beta_r'' = A_r + B_{r-1} + Y_r + X_r \dots (63)$$

Which is the general form of the Theorem of three moments. It expresses the relation between the bending moments over three consecutive supports in terms of the loads, spans, E. I and h. An equation of the form of (63) can be written for every support, and, as, if the girder rests on the supports, the moment at the first and last support is zero, we shall have as many equations as there are unknown quantities or bending moments, and hence we can determine their values.

Let there be s spans in the continuous girder represented in Fig. 1, and let the rth span alone be loaded, then by the three moment theorem, the equation for each support is as follows:

Multiplying the first equation of (64) by c_2 , the second by c_3 , the r^{th} by c_r , etc., we obtain, after reduction,

Now, supposing it is desired to determine the value of M_s , it is only necessary to impose such conditions upon the multiplier e that all terms shall reduce to zero, excepting those containing M_s . Evidently, the co-efficients must separately equal zero, or

And, for M_s , we have, at once,

$$M_{s} = \frac{(Y_{r} + X_{r}) c_{r} + A_{r} c_{r} + B_{r} c_{r+\ell}}{(c_{s-\ell} \beta_{s-\ell}'' + 2 c_{s} \beta_{s}') = -c_{s+\ell} \beta_{s}} . \qquad (67)$$

In a similar manner, multiplying the last equation of (64) by d_2 , the last but one by d_2 , and so on, we obtain

$$\begin{cases}
2 d_{z} \beta'_{s-i} + d_{s} \beta''_{s-i} = 0 \\
d_{z} \beta_{s-i} + 2 d_{s} \beta'_{s-i} + d_{4} \beta''_{s-2} = 0 \\
d_{s} \beta_{s-2} + 2 d_{4} \beta'_{s-2} + d_{5} \beta''_{s-3} = 0 \\
\times \times \times \times \times \times \\
d_{m-i} \beta_{s-m+2} + 2 d_{m} \beta'_{s-m+2} + d_{m+i} \beta''_{s-m+i} = 0
\end{cases}$$
(68)

And

$$M_{2} = \frac{(Y_{r} + X_{r}) d_{s-r+2} + A_{r} d_{s-r+2} + B_{r} d_{s-r+i}}{(2 \beta_{2}' d_{s} + d_{s-i} \beta_{2}) = -d_{s+i} \beta_{i}'} \dots (69)$$

From (65), we see that $M_s = -\frac{2}{\beta_2'}\frac{\beta_2'}{\beta_2''}M_z$, from (66), $c_s = -\frac{2}{2}\frac{c_z}{\beta_2'}$; hence, assuming $c_i = 0$ and $c_z = 1$, we have $M_s = c_s$ $\frac{\beta_2}{\beta_2'}M_z$, and in a similar manner we find that $M_4 = c_4$ $\frac{\beta_2}{\beta_2'}\frac{\beta_3}{\beta_2''}$ M_s , $M_s = c_s$ $\frac{\beta_2}{\beta_2''}\frac{\beta_3}{\beta_2''}\frac{\beta_4}{\beta_2''}M_z$, or in general for any support on the left of the loaded span or spans.

In a like manner, we obtain for any support on the right of the loaded span or spans,

$$m > r$$

$$M_m = d_{s-m+2} \frac{\beta_{s-1}'' \beta_{s-2}'' \cdot \beta_m''}{\beta_{s-1}' \beta_{s-2} \cdot \beta_m} M_s \cdot \dots$$

$$(71)$$

Substituting (69) in (70), and (67) in (71), we obtain m < r+1

$$M_{m} = -\frac{c_{m} \beta_{2} \beta_{3} \beta_{3} \beta_{3} \beta_{m-1}}{(d_{s+1} \beta_{1}^{"}) \beta_{2}^{"} \beta_{2}^{"} \beta_{2}^{"}} \left\{ (A_{r} + Y_{r} + X_{r}) d_{s-r+2} B_{r} d_{s-r+1} \right\} (72)$$

$$M_{m} = -\frac{d_{s-m+2} \beta_{s-i}^{"} \beta_{s-i}^{"} \beta_{s-2}^{"} \dots \beta_{m}^{"}}{(c_{s+i} \beta_{s}) \beta_{s-i} \beta_{s-2} \dots \beta_{m}} \left\{ (A_{r} + Y_{r} + X_{r}) c_{r} + B_{r} c_{r+i} \right\} \dots (73)$$

In (72), as m must always be less than r+1, r can have values from s to m.

In (73), as m must always be greater than r, r can have values from 1 to m-1.

Hence, adding (72) and (73), and substituting $X'_r + X''_{r-1}$ for X_r , we have

$$M_{m} = \frac{-c_{m}\beta_{2}\beta_{3}...\beta_{m-1}}{(d_{s+1}\beta'_{1})\beta''_{2}\beta''_{3}...\beta''_{m-1}} \frac{\Sigma}{r=s} \left\{ (A_{r}+Y_{r}+X'_{r}+X''_{r-1})d_{s-r+2}+B_{r}d_{s-r+1} \right\}$$

$$+\frac{-d_{s-m+2}}{(c_{s+1}\beta_{s})}\frac{\beta_{s-1}''\beta_{s-2}''}{\beta_{s-1}\beta_{s-2}\ldots\beta_{m}''}\sum_{r=m-1}^{r=1}\left\{\left(A_{r}+Y_{r}+X_{r}'+X_{r-1}''\right)c_{r}+B_{1}c_{r+1}\right\}\left(A\right)$$

From (A), we can obtain the bending moment over any support m of a continuous girder of any number of spans s, of any lengths as $l_1, l_2 \ldots l_s$, supports at any levels, the moment of inertia I constant or variable, the modulus of elasticity E being constant, and the loads being placed at pleasure.

Note that $\beta_2 \leq \beta_{m-1}^n$, $\beta_2 \leq \beta_{m-1}$, $\beta_{s-1}^n \geq \beta_m^n$, and $\beta_{s-1} \geq \beta_m$.

$k = \frac{\epsilon t}{l}$	$k-k^{3}$		$k = \frac{a}{l}$	$k-k^3$		$k = \frac{\alpha}{1}$	$k-k^3$	
.001 2 3 4 5 6 7 8	C00 999 999 001 999 992 002 999 973 003 999 936 004 999 875 005 999 784 006 999 657 007 999 488 C08 999 271	.999 8 7 6 5 4 3 2	.060 61 62 63 64 65 66 67 68 69	059 784 C00 060 773 019 061 761 672 062 749 953 063 737 856 064 725 375 065 712 504 066 699 237 067 685 568 068 671 491	.940 39 38 37 36 35 34 33 32 31	.120 21 22 23 24 25 26 27 28 29	118 272 000 119 228 439 120 184 152 121 139 133 122 093 376 123 046 875 123 999 624 124 951 617 125 902 848 126 853 311	.880 79 78 77 76 75 74 73 72 71
.010 11 12 13 14 15 16 17 18	009 989 000 010 998 669 011 998 272 012 997 803 013 997 256 014 996 625 015 995 904 016 995 087 017 994 168 018 993 141	.990 89 88 87 86 85 84 83 82 81	.070 71 72 73 74 75 76 77 78 79	069 657 000 070 642 089 071 626 752 072 610 983 073 594 776 074 578 125 075 561 024 076 543 467 077 525 448 078 506 961	.930 29 28 27 26 25 24 23 22 21	31 32 33 34 35 36 37 38 39	127 803 000 128 751 909 129 700 032 130 647 363 131 593 896 132 539 625 133 484 544 134 428 647 135 371 928 136 314 381	.870 69 68 67 66 65 64 63 62 61
.020 21 22 23 24 25 26 27 28 29	019 992 000 020 990 789 021 989 352 022 987 833 023 986 176 024 984 375 025 982 424 026 980 317 027 978 048 028 975 611	.980 79 78 77 76 75 74 73 72 71	.080 81 82 83 84 85 86 87 88 89	079 488 000 080 468 559 081 448 632 082 428 213 083 407 296 084 385 875 085 363 944 086 341 497 087 318 528 088 295 031	.920 19 18 17 16 15 14 13 12	.140 41 42 43 44 45 46 47 48 49	157 256 000 138 196 779 139 136 712 140 075 793 141 014 016 141 951 375 142 887 861 143 823 477 144 758 208 145 692 051	.860 59 58 57 56 55 54 53 52 51
.030 31 32 33 34 35 36 37 38 39	029 973 000 030 970 209 031 967 232 032 964 063 033 960 696 034 957 125 035 953 344 036 949 347 037 945 128 038 940 681	.970 69 68 67 66 65 64 63 62 61	.090 91 92 93 94 95 96 97 98 99	089 271 000 090 246 429 091 221 312 092 195 63 093 169 416 094 142 625 095 115 264 096 087 327 097 058 808 098 029 701	.910 9 8 - 7 6 5 4 3 2 1	.150 51 52 53 54 55 56 57 58 59	146 625 000 147 557 049 148 488 192 149 418 423 150 347 736 151 276 125 152 203 584 153 130 107 154 055 688 154 920 321	.850 49 48 47 46 45 44 43 42 41
.040 41 42 43 44 45 46 47 48 49	039 936 000 040 931 079 041 925 912 042 920 493 043 914 816 044 908 875 045 902 664 046 896 177 047 889 408 048 882 351	.960 59 58 57 56 55 54 53 52 51	.100 1 2 3 4 5 6 7 8 9	099 000 000 099 969 699 100 938 792 101 907 273 102 875 136 103 842 375 104 808 984 105 774 957 106 740 288 107 704 971	900 899 98 97 96 95 94 93 92 91	.160 61 62 63 64 65 66 67 68 69	155 904 000 156 826 719 157 748 472 158 669 258 159 589 056 160 507 875 161 425 704 162 342 587 163 258 368 164 173 191	.810 39 38 37 36 35 34 33 32 31
.050 51 52 53 54 55 56 57 58 59	049 875 000 050 867 349 051 859 392 052 851 123 053 842 536 054 813 625 055 824 884 056 814 807 057 804 888 058 794 621	.950 49 48 47 46 45 44 43 42 41	.110 11 12 13 14 15 16 17 18 19	108 669 000 109 632 369 110 595 072 111 557 103 112 518 456 113 479 125 114 439 104 115 398 387 116 356 968 117 314 841	.890 89 88 87 86 85 84 83 82 81	.170 71 72 73 74 75 76 77 78 79	165 087 000 165 999 789 166 911 552 167 822 283 168 731 976 169 640 625 170 548 224 171 454 767 172 360 248 173 264 661	.830 29 28 27 26 25 24 23 22 21
		$k = \frac{a}{l}$		$2 k-3 k^3+k^3$				$k = \frac{a}{l}$

1	1	1	1		1	1 1		
$k = \frac{\alpha}{l}$	$k-k^3$			$k-k^3$		$k = \frac{a}{l}$	$k-k^3$	
.180 81 82 83 84 85 86 87 88	174 168 009 175 070 259 175 971 482 176 871 513 177 770 496 178 668 375 179 595 144 180 460 797 181 355 328 182 248 731	.82) 19 18 17 16 15 14 13 12	1	226 176 000 227 002 479 227 827 512 228 651 093 229 473 216 230 293 875 231 113 061 231 930 777 232 747 008 233 561 751	.760 59 58 57 56 55 54 53 52 51	.300 1 2 3 4 5 6 7 8 9	273 000 000 273 729 099 274 456 392 275 181 873 275 905 536 276 627 375 277 347 384 278 065 557 278 781 888 279 496 371	.700 .699 .98 .97 .96 .95 .91 .93 .92
.190 91 92 93 94 95 96 97 98	193 141 000 184 032 129 184 922 112 185 810 943 186 698 616 187 585 125 183 470 464 189 354 627 190 237 608 191 119 401	.810_9 8 77 6 5 4 3 2 1	.250 51 52 +3 54 55 56 57 58 59	234 375 000 235 186 749 235 996 992 236 805 723 237 612 936 238 418 6.5 239 222 78 240 025 407 240 826 488 241 626 021	.750 49 48 47 46 45 41 43 42 41	.310 11 12 13 14 15 16 17 18	280 209 000 280 919 769 281 628 672 282 335 703 283 040 856 283 744 125 284 445 504 285 144 987 285 842 568 286 538 241	.690 89 88 87 86 85 84 83 82 81
.200 1 2 3 4 5 6 7 8 9	192 000 000 192 879 399 193 757 592 194 634 573 195 510 336 196 384 875 197 258 184 198 130 257 199 001 088 199 870 671	.8°0 .799 98 97 96 95 94 93 92 91	.260 61 62 63 64 65 66 67 68	742 424 000 243 220 419 244 015 272 244 808 558 245 600 256 246 390 375 247 178 904 247 965 837 248 751 168 249 534 891	.740 39 38 37 36 35 34 33 32 31	.320 21 22 23 24 25 26 27 28 29	287 232 000 287 923 839 288 613 752 289 301 733 289 987 776 290 671 875 291 354 024 292 034 217 292 712 448 293 388 711	.680 79 78 77 76 75 74 73 72 71
.210 11 12 13 14 15 16 17 18 19	200 739 000 201 606 069 202 471 872 203 336 403 204 199 656 205 061 625 205 922 304 206 781 687 207 639 768 208 496 541	.790 89 88 87 86 85 84 83 82 81	.270 71 72 73 74 75 76 77 78 79	250 317 000 251 097 489 251 876 352 252 653 583 253 429 176 254 203 125 254 975 424 255 746 067 256 515 048 257 282 361	.780 29 28 27 26 25 24 23 22 21	.330 31 32 33 34 35 36 37 38 39	291 063 000 294 735 309 295 405 682 296 0.3 963 296 740 296 297 404 625 298 066 944 298 727 247 299 385 528 300 041 781	.670 69 68 67 66 65 64 63 62 61
.220 21 22 23 24 25 26 27 28 29	209 352 000 210 206 139 211 058 952 211 910 433 212 760 576 213 609 375 214 456 824 215 302 917 216 147 648 216 991 011	.780 79 78 77 76 75 74 73 72 71	.280 81 82 83 84 85 86 87 88 89	258 048 000 258 811 959 259 574 232 260 334 813 261 093 696 261 850 875 262 606 344 263 360 097 264 112 128 264 862 431	.720 19 18 17 16 15 14 13 12 11	.340 41 42 43 44 45 46 47 48 49	30) 696 000 301 348 179 301 998 312 302 646 393 303 292 416 303 936 375 304 578 264 505 218 077 305 855 808 306 491 451	.660 59 58 57 56 55 54 53 52 51
.230 31 32 33 34 35 36 37 38 39	217 833 000 218 673 609 219 512 832 220 350 663 221 187 096 222 022 125 222 855 747 223 687 947 224 518 728 225 348 081	.770 69 68 67 66 65 64 63 62 61	.290 91 92 93 94 95 96 97 98 99	265 611 0(0 266 357 829 267 102 912 267 846 243 268 587 816 269 327 625 270 065 664 270 801 927 271 536 408 272 269 101	710 9 8 7 6 5 4 3 2	.350 51 52 53 54 55 56 57 58 59	307 125 000 307 756 449 308 385 792 309 013 023 309 638 136 310 261 125 310 881 934 311 500 707 312 117 288 312 731 721	.650 49 48 47 46 45 44 -43 42 41
	$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$		2 k-3 k ² +k ³	$k = \frac{a}{l}$		$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$

$k = \frac{a}{l}$	$k-k^3$		$k = \frac{a}{l}$	$k-k^3$		$k = \frac{\alpha}{l}$	$k-k^3$	
.360 61 62 63 64 65 66 67 68 69	313 344 000 313 954 119 314 562 072 315 167 853 315 771 456 316 372 875 316 972 104 317 569 137 318 163 968 318 756 591	.649 39 38 38 37 36 35 34 33 32 31	.420 21 22 23 24 25 26 27 28 29	345 912 +60 346 381 539 346 848 552 347 313 033 347 774 976 348 234 375 348 691 224 349 145 517 349 597 248 350 046 411	.580 79 78 77 76 75 74 78 72 71	.4 0 81 82 83 84 85 86 87 88 89	408 C0) 715 359 370 019 832 321 413 620 096 915 875 371 208 744 498 697 785 728 372 069 831	.520 19 18 17 16 15 14 13 12 11
.370 71 72 73 74 75 76 77 78 79	319 347 000 319 935 189 320 521 152 321 104 883 321 686 376 322 265 625 322 842 624 323 417 367 323 989 848 324 560 061	.630 29 28 27 26 25 - 24 23 22 21	.430 31 32 33 34 35 36 37 38 39	550 493 000 350 937 009 351 378 432 351 817 263 352 253 496 352 687 125 353 118 144 353 546 547 353 972 328 354 395 481	69 68 67 66 65 64 63 62 61	,450 91 92 93 94 95 96 97 98	351 000 629 229 904 512 373 176 843 446 216 712 625 976 064 374 236 527 494 008 748 501	.510 9 8 7 6 5 4 3 2
.380 81 82 83 84 85 86 87 88	325 128 000 325 693 659 326 257 032 326 818 113 527 376 896 327 933 375 323 487 544 329 039 397 329 588 928 33) 136 131	.620 19 18 17 16 15 14 13 12	43 44 45 46 47 48	354 816 000 355 233 879 355 649 112 356 661 693 356 471 616 356 878 875 357 283 464 357 685 377 358 084 608 358 481 151	.560 59 58 57 56 £5 54 53 52 51		375 248 499 498 992 736 473 975 936 376 212 375 445 784 676 157 903 488 377 127 771	.500 .499 .98 .97 .96 .95 .94 .93 .92 .91
.390 91 92 93 94 95 96 97 98	3 0 681 000 331 223 529 331 763 712 332 301 543 332 837 016 333 370 125 333 900 864 334 429 227 334 955 208 335 478 801		.450 51 52 53 54 55 56 57 58 59		.559 49 48 47 46 45 44 43 42 41	.510 11 12 13 14 15 16 17 18 19	349 000 567 169 782 272 994 3.3 378 203 256 409 125 611 904 811 587 379 008 168 201 641	.490 89 88 87 86 85 84 83 82 81
.400 1 2 3 4 5 6 7 8 9	336 000 000 336 518 799 337 035 192 337 549 173 338 060 736 338 569 875 339 076 584 339 580 857 340 082 688 340 582 071	.600 .599 98 97 96 95 94 93 92 91	.460 61 62 63 64 65 66 67 68 69	664 000 863 027 819 363 388 872 747 153 364 102 656 455 375 805 304 365 152 437 496 768 838 291	.540 39 38 37 36 35 34 33 32 31	.520 21 22 23 24 25 26 27 28 29		.480 79 78 77 76 75 74 73 72 71
.410 11 12 13 14 15 16 17 18 19	341 079 000 341 573 469 342 065 472 342 555 003 343 042 056 343 526 (25 344 008 704 344 488 287 344 965 368 345 439 941	.590 89 88 87	.470 71 72 73 74 75 76 77 78 79	366 177 000 512 889 845 952 367 176 183 503 576 828 125 . 368 149 824 468 667 784 648 369 097 761	.530 29 28 27 26 25 24 23 22 21	.£30 31 32 33 34 35 36 37 38 39	381 123 000 278 709 431 232 580 563 726 696 869 625 382 009 344 145 847 279 128 409 181	.470 69 63 67 66 65 64 63 62 61
	$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$		$2k-3k^2+k^3$			$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$

$k = \frac{\alpha}{l}$	$k-k^3$		$ k=\frac{a}{l} $	$k-k^{s}$		$k = \frac{\alpha}{l}$	kk3	
.540 41 42 43 44 45 46 47 48	536 0d4 659 579 779 912 886 993 383 010 816 121 375 228 664 332 677 433 408 530 851	.460 59 58 57 56 55 54 53 52 51	.600 1 2 3 4 5 6 7 8 9	384 000 000 383 918 199 832 792 743 773 651 136 554 875 454 984 351 457 244 288 133 471	.400 .399 98 97 96 95 94 93 92 91	660 61 62 63 64 65 66 67 68 69	504 + 00 195 - 219 371 - 882 - 472 565 - 753 245 - 056 370 - 920 - 375 591 - 704 259 - 037 369 - 922 - 368 581 - 691	340 39 38 37 36 35 34 33 32 81
.550 51 52 53 54 55 56 57 58 59	383 625 000 715 849 803 3942 887 623 968 536 384 046 125 120 384 191 307 250 888 323 121	.450 49 48 47 46 45 44 43 42 41	.610 11 12 13 14 15 16 17 18	019 000 382 900 869 779 072 653 603 524 456 391 625 255 104 114 887 381 970 968 823 341	.390 89 88 87 86 81 84 83 82 81	.670 71 72 73 74 75 76 77 78 79	237 000 368 888 289 535 552 178 783 367 817 976 453 125 084 224 366 711 267 334 248 365 953 161	.330 29 28 27 26 25 24 23 22 21
.560 61 62 63 64 65 66 67 68 69	384 000 441 519 495 672 546 453 593 836 637 875 678 504 715 787 749 568 779 991		.620 21 22 23 24 25 26 27 28 29	672 000 516 939 358 152 195 633 029 376 380 859 375 685 624 508 117 326 848 141 811	.380 79 78 77 76 75 74 73 72 71	.680 81 82 83 84 85 86 87 88	568 000 178 759 361 785 482 588 013 363 986 496 580 875 171 144 362 757 297 339 328 361 917 231	.320 19 18 17 16 15 14 13 12
570 71 72 73 74 75 76 77 78 79	807 000 830 589 850 752 867 483 880 776 890 625 897 024 899 967 899 448 895 461	.430 29 28 27 26 25 24 23 22 21	630 31 32 33 34 35 36 37 38 39	379 953 000 760 409 564 032 363 863 363 863 159 896 378 952 125 740 £44 525 147 305 928 082 881	.370 63 68 67 66 65 64 63 62 61	.690 91 92 93 94 95 96 97 98 99	491 000 060 629 360 626 112 187 443 359 744 616 297 625 358 846 464 391 127 357 931 608 467 901	.310 9 8 7 6 5 4 3 2
.580 81 82 83 84 85 86 87 88 89	888 00 877 059 862 632 844 713 823 296 798 375 769 944 737 997 702 528 663 531	.420 19 18 17 16 15 14 13 12	.640 41 42 43 44 45 46 47 48 49	377 856 000 625 279 390 712 152 293 376 910 016 663 875 413 864 159 977 375 902 208 640 551	.360 59 58 57 56 55 54 53 52 51	.700 1 2 3 4 5 6 7 8 9	357 009 000 356 527 899 051 592 355 571 073 086 336 354 597 375 354 104 184 353 606 757 105 088 352 599 171	.307 .299 .98 .97 .96 .95 .94 .93 .92 .91
.590 91 92 93 94 95 96 97 98 99	621 000 574 929 525 312 472 143 415 416 355 125 291 264 223 827 152 808 078 201	.410 9 8 7 6 5 4 3 2	.6 0 51 52 53 54 55 56 57 58 59	375 575 000 105 549 374 832 192 554 923 273 736 373 988 625 699 584 406 607 109 688 372 808 821	.350 49 48 47 46 45 44 43 42 41	.710 11 12 13 14 15 16 17 18	(89 000 351 574 569 055 872 350 532 903 005 656 349 474 125 348 938 304 398 187 347 853 768 305 041	.290 89 88 87 86 85 84 83 82 81
	$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$		$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$		$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$

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$k = \frac{\alpha}{l}$	$k-k^3$		$k = \frac{l}{l}$	$k-k^{s}$		$k = \frac{\alpha}{l}$	$k-k^3$	
.720 21 22 23 24 25 26 27 28 29	346 752 000 194 639 345 632 952 666 933 344 496 576 343 921 875 343 342 824 342 759 417 171 648 341 579 511	.280 79 78 77 76 75 74 73 72 71	.780 81 82 83 84 85 86 87 88 89	305 448 000 304 620 459 303 788 232 302 951 313 302 109 696 301 263 375 300 412 344 299 556 597 293 696 128 297 830 931	.220 19 18 17 16 15 14 13 12	.840 41 42 43 44 45 46 47 48 49	247 296 000 246 176 679 245 052 312 243 922 893 242 788 416 241 648 875 240 504 264 239 354 577 238 199 808 237 039 951	.160 59 58 57 56 55 54 53 52 51
.730 31 32 33 34 35 36 37 38 39	340 983 000 382 109 389 776 882 167 163 388 553 096 337 934 625 311 744 336 684 447 052 728 335 416 581	.270 69 68 67 66 65 64 63 62 61	.790 91 92 93 94 95 96 97 98	296 961 000 296 086 329 295 206 912 294 322 743 293 433 816 292 540 125 291 641 664 290 738 427 289 830 408 288 917 601	.210 9 8 7 6 5 4 3 2 1	.850 51 52 53 54 55 56 57 58 59	235 875 000 234 704 919 233 529 792 232 349 523 231 164 136 229 973 625 228 777 984 227 577 207 226 371 288 225 160 221	.150 49 48 47 46 45 44 43 42 41
.740 41 42 43 44 45 46 47 48 49	334 776 000 130 979 333 481 512 332 827 593 169 216 331 506 375 330 839 064 167 277 329 491 008 328 810 251	.260 59 58 57 56 55 54 53 52 51	.800 1 2 3 4 5 6 7 8 9	288 000 000 287 077 599 286 150 392 285 218 373 284 281 536 283 339 875 282 393 384 281 442 057 280 485 888 279 524 871	.200 .199 .98 .97 .96 .95 .94 .93 .92	.860 61 62 63 64 65 66 67 68 69	223 944 000 222 722 619 221 496 072 220 264 353 219 027 456 217 785 375 216 538 104 215 285 637 214 027 968 212 765 091	.140 39 38 37 36 35 34 33 32 31
.750 51 52 53 54 55 56 57 58 59	328 125 000 327 435 249 326 740 992 326 042 223 325 338 936 324 631 125 328 918 784 323 201 907 322 480 488 321 754 521	.250 49 48 47 46 45 44 43 42 41	.810 11 12 13 14 15 16 17 18 19	278 559 000 277 588 269 276 612 672 275 632 203 274 646 856 273 656 625 272 661 504 271 661 487 270 656 568 269 646 741	.190 89 88 87 86 85 84 83 82 81	.870 71 72 73 74 75 76 77 78 79	211 497 000 210 223 689 208 945 152 207 661 383 206 372 376 205 078 125 203 778 624 202 473 867 201 163 848 199 848 561	.130 29 28 27 26 25 24 23 22 21
.760 61 62 63 64 65 66 67 68 69	321 024 000 320 288 919 319 549 272 318 805 053 318 056 256 317 302 875 316 544 904 315 782 337 315 015 168 314 243 391	.240 39 38 37 36 35 34 33 32 31	.820 21 22 23 24 25 26 27 28 29	268 632 000 267 612 339 266 587 752 265 558 233 264 523 776 263 484 375 262 440 024 261 390 717 260 336 448 259 277 211	.180 79 78 77 76 75 74 73 72 71	.880 81 82 83 84 85 86 87 88 89	198 528 000 197 202 159 195 871 032 194 534 613 193 192 896 191 845 875 190 493 544 189 135 897 187 772 928 186 404 631	.120 19 18 17 16 15 14 13 12 11
.770 71 72 73 74 75 76 77 78 79	313 467 000 312 685 989 311 900 352 311 110 083 310 315 176 309 515 625 308 711 424 307 902 567 307 089 048 306 270 861	.230 29 28 27 26 25 24 23 22 21	31 32 33 34 35 36	258 213 000 257 143 809 256 069 632 254 990 463 253 906 296 252 817 125 251 722 944 250 623 747 249 519 528 248 410 281	.170 69 68 67 66 65 64 63 62 61	.890 91 92 93 94 95 96 97 98	185 031 000 183 652 029 182 267 712 180 878 043 179 483 016 178 082 625 176 676 864 175 265 727 173 849 208 172 427 301	.110 9 8 7 6 5 4 3 2 1
	$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$		2 k-3 k³+k³	$k = \frac{a}{l}$		$2 k-3 k^2+k^3$	$k = \frac{a}{l}$

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$k = \frac{a}{l}$	$k-k^3$		$k = \frac{a}{l}$	$k-k^3$		$k = \frac{a}{l}$	$k-k^3$	
.900 1 2 3 4 5 6 7 8	171 000 000 169 5e7 299 168 129 192 166 685 673 165 236 736 163 782 375 162 322 584 160 857 357 159 386 688 157 910 571	.100 99 98 . 97 96 95 94 93 92 91	.940 41 42 43 44 45 46 47 48 49	109 416 000 107 762 379 106 103 112 104 438 193 102 767 616 101 091 375 099 409 464 097 721 87 096 028 608 094 329 651	60 59 58 57 56 55 54 53 52 51	.980 81 82 83 84 85 86 87 88 89	033 808 000 036 923 859 035 033 832 033 137 913 031 236 098 029 328 375 027 414 744 025 495 197 023 569 728 021 638 331	20 19 18 17 16 15 14 13 12 11
.910 11 12 13 14 15 16 17 18 19	156 429 000 154 941 969 153 449 472 151 951 503 150 448 056 148 939 125 147 424 704 145 904 787 144 379 368 142 848 441	90 89 88 87 86 85 84 83 82 81	.950 51 52 53 54 55 56 57 58 59	092 625 000 090 914 649 089 198 592 087 476 823 085 749 336 084 016 125 082 277 184 (80 532 507 078 782 088 077 025 921	50 49 48 47 46 45 41 43 42 41	.990 91 92 93 94 95 96 97 98 99	019 701 000 017 757 729 015 808 512 013 853 343 011 892 216 009 925 125 007 952 064 005 973 027 003 995 0(8 001 997 001	10 9 8 7 6 5 4 3 2
.920 21 22 23 24 25 26 27 28 29	141 312 000 139 770 039 138 222 552 136 669 533 135 110 976 133 *46 875 131 977 224 130 402 017 128 821 248 127 234 911	80 79 78 77 76 75 74 73 72 71	.96) 61 62 63 64 65 66 67 68 69	075 264 000 073 496 319 071 722 872 069 943 653 068 158 656 066 367 875 064 571 304 062 768 937 060 960 768 059 146 791	40 39 38 37 36 35 34 33 32 31		2 k-3 k²+k³	$k = \frac{a}{l}$
.930 31 32 33 34 35 36 37 38 39	125 643 000 124 045 509 122 442 432 120 838 763 119 219 496 117 599 625 115 974 144 114 343 047 112 706 328 111 (63 981	70 69 68 67 66 65 64 63 62 61	.970 71 72 73 74 75 76 77 78 79	057 327 000 055 501 389 053 669 952 051 832 683 049 989 576 048 140 625 046 285 824 044 425 167 042 558 648 040 686 261	30 29 28 27 26 25 24 23 22 21			
	$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$		$2 k - 3 k^2 + k^3$	$k = \frac{a}{l}$			



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^{*} Variable Moment of Inertia.

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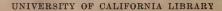












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